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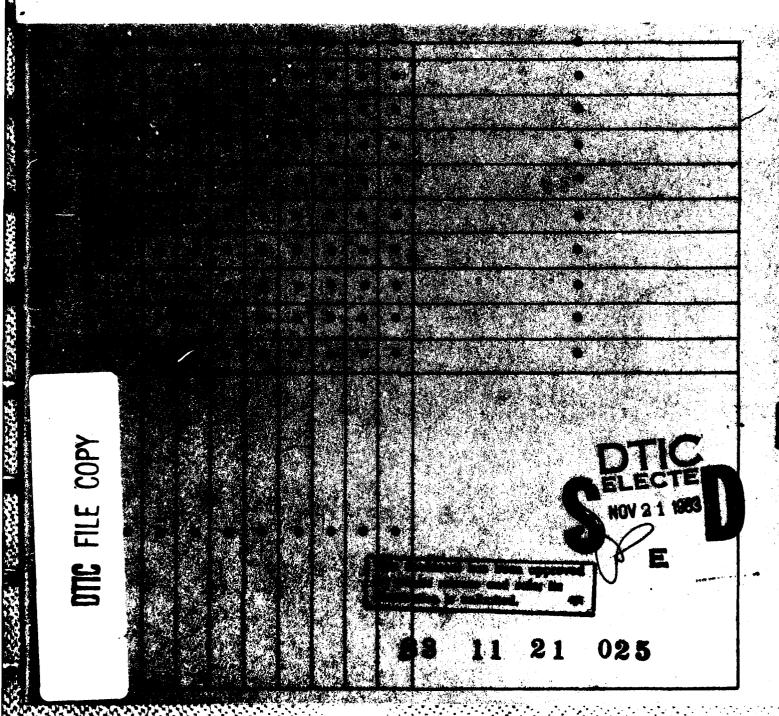
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US Army Corps of Engineers

Cold Regions Research & Engineering Laboratory

Computer models for two-dimensional transient heat conduction



For conversion of SI metric units to U.S./British customary units of measurement consult ASTM Standard E380, Metric Practice Guide, published by the American Society for Testing and Materals, 1916 Race St., Philadelphia, Pa. 19103.

Cover: Finite difference grid of nodes for the semi-infinite corner problem.





Computer models for two-dimensional transient heat conduction

Mary Remley Albert

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This paper documents the development and verification of two finite difference models that solve the general two-						
dimensional form of the heat conduction equation, using the alternative-direction implicit method. Both can handle						
convective, constant flux, specified temperature and semi-infinite boundaries. The conducting medium may be com-						
convective, constant flux, specified temperature and semi-infinite boundaries. The conducting medium may be composed of many materials. The first program, ADI, solves for the case where no change of state occurs. ADIPC solves for the case where a freeze/thaw change of phase may occur, using the apparent heat capacity method. Both models						

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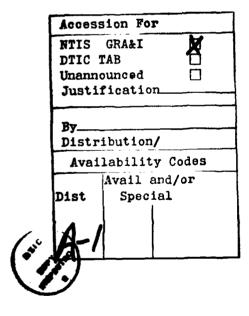
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PREFACE

This report was prepared by Mary Remley Albert, Mathematician, of the Applied Research Branch, Experimental Engineering Division, U.S. Army Cold Regions Research and Engineering Laboratory. The work was sponsored by DA Project 4A762730AT42, Design, Construction and Operations Technology for Cold Regions, Technical Area D, Cold Regions Design and Construction, Work Unit 017, Heat Distribution Systems in Cold Regions.

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COMPUTER MODELS FOR TWO-DIMENSIONAL TRANSIENT HEAT CONDUCTION

Mary Remley Albert

INTRODUCTION

Most major Army installations are heated with central heat distribution systems. The thermal regimes around buried distribution systems are of interest for several reasons, not the least of which is the estimated millions of dollars lost annually from these systems (Phetteplace et al. 1981). Also, when replacing damaged insulation or installing insulation in a new system, it is desirable to know what an optimum balance is between the initial cost of insulation and the continued operating cost of heat losses. In addition, freezing and thawing of the ground around a buried distribution system could damage it, through loss of support and settlement, for example.

Analysis of the thermal regime involves the solution of the heat conduction equation for situations with complicated geometries and a variety of boundary conditions. Usually, these partial differential equations cannot be solved analytically; we must resort to numerical methods. One long-established method is that of finite differences, a relatively straightforward numerical method used successfully to solve a variety of differential equations. A finite difference computer program, set up in a general form to solve the heat conduction equation under a variety of geometries and boundary conditions, provides a powerful tool with which the engineer can assess problems in conductive heat transfer.

The objective of this paper is to document the development and verification of two general two-dimensional finite difference computer programs that were written to model time-varying heat conduction in a medium composed of many materials. The first program, ADI, solves for the general case in which no change of state occurs in the conducting medium. The second program, ADIPC, is an adaptation of ADI that includes the effects of phase change in the conducting medium. The programs are able to handle convective, constant flux, specified temperature and semi-infinite boundaries. Material properties such as thermal conductivity, density and specific heat are allowed to vary with time or temperature. The programs are relatively easy to use, that is, they are easily set up for new conduction problems by those who have a background knowledge of FORTRAN.

The computer programs were written and programmed in FORTRAN by the author on CRREL's PRIME 400 computer.

FINITE DIFFERENCES APPLIED TO HEAT TRANSFER

Finite differences are commonly used in the numerical solution of partial differential equations. The procedure involves the replacement of differentials by differences, and is best illustrated by construction of the so-called finite difference grid. Let us examine a grid that is set up in two dimensions to model heat conduction (Fig. 1). The two dimensions in this case represent spatial independent variables, x and y. Each point of the grid is called a node and represents the area enclosed by the square around it (with unit depth). T(x, y) represents the temperature of node x, y. For each node, the temperature and material properties are assumed uniform for the region it represents. By specifying the initial temperatures at each of the nodes, the nodal thermal properties and the boundary conditions, we can solve the heat conduction equation to determine the temperatures of the nodes at later times.

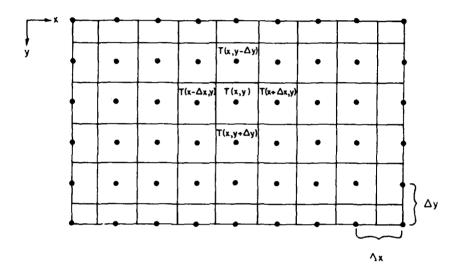


Figure 1. A simple finite difference grid.

Heat conduction equation

The equation governing transient heat conduction in two dimensions is

$$\frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) = \rho C_{p} \frac{\partial T}{\partial t} \tag{1}$$

where x and y = spatial variables

T = temperature

t = time

k = thermal conductivity (this may vary over the spatial and time domains)

 ρ = density

 C_n = specific heat.

The nonlinearities in this partial differential equation that arise from the inclusion of the phase change condition will be discussed in the *Phase Change* section.

Partial differential equations may be expressed in finite difference form either by a Taylor's series expansion about a point or by physical considerations, such as a heat flow balance in the case of the heat conduction problem. For problems involving variable thermal properties or complicated boundary conditions, the heat balance approach is the simplest.

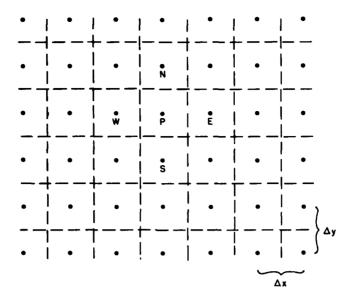


Figure 2. A node on the interior of the grid.

Consider the control volume represented by a node on the interior of the grid, illustrated in Figure 2. The heat flow Q into the control volume from any adjacent node may be calculated as follows:

$$Q = \frac{k}{L} A \Delta T \tag{2}$$

where

k = thermal conductivity

L =distance over which the heat flow occurs

A = area (per unit depth) perpendicular to the direction of heat flow

 ΔT = change in temperature between the two nodes.

Examine the heat flow from node N to node P. The conductance k/L is the reciprocal of the resistance between the two nodes. The resistance between the two nodes $R_{\rm NP}$ is simply the sum of the individual resistances,

$$R_{NP} = R_N + R_p$$
.

Each resistance is the distance divided by the material conductivity,

$$R_{\rm NP} = \frac{\Delta y/2}{k_{\rm N}} + \frac{\Delta y/2}{k_{\rm p}}$$

then

$$\frac{k}{L} = \frac{1}{R_{\rm NP}} = \frac{2 k_{\rm N} k_{\rm P}}{\Delta y (k_{\rm N} + k_{\rm p})} \ . \tag{3}$$

In the computer program it is assumed that, for nodes not on a boundary of the grid, Δy and Δx are uniform and equal throughout the grid. Let the nodal spacing $\Delta x = \Delta y = \Delta s$. Following the form of eq 2, we find that the heat flows into the region associated with node P from nodes N, W, S and E may be calculated as follows:

$$Q_{NP} = \frac{2 k_N k_P}{\Delta s (k_N + k_P)} \quad (\Delta s) (T_N - T_P) \tag{4}$$

$$Q_{WP} = \frac{2k_W k_P}{\Delta s(k_W + k_P)} (\Delta s) (T_W - T_P)$$
 (5)

$$Q_{SP} = \frac{2k_S k_P}{\Delta s (k_S + k_P)} (\Delta s) (T_S - T_P)$$
(6)

$$Q_{\rm EP} = \frac{2k_{\rm E}k_{\rm P}}{\Delta s(k_{\rm E} + k_{\rm P})} (\Delta s) (T_{\rm E} - T_{\rm P}) \tag{7}$$

Note that the area perpendicular to the direction of heat flow is taken per unit depth, that is, $A = \Delta s \cdot 1$.

The sum of the heat flows into a node is responsible for the temperature change of the node:

$$Q_{NP} + Q_{WP} + Q_{SP} + Q_{EP} = V_{P} \cdot \frac{1}{\Delta t} \left(\rho' C_{P}' T_{P}' - \rho C_{P} T_{P} \right)$$
 (8)

where $V_{\rm p}$ is the control volume for unit depth ($V_{\rm p} = (\Delta s)^2$ in this case) and the primed variables represent time $t + \Delta t$.

Now substitute eq 4-7 into eq 8 to get the finite difference form of eq 1 for a node not on the boundary of the grid:

$$k_{\rm NP}(T_{\rm N} - T_{\rm P}) + k_{\rm WP} (T_{\rm W} - T_{\rm P}) + k_{\rm SP} (T_{\rm s} - T_{\rm P}) + k_{\rm E} (T_{\rm E} - T_{\rm P}) = \frac{(\Delta s)^2}{\Delta t} (\rho' C_{\rm P}' T_{\rm P}' - \rho C_{\rm P} T_{\rm P}).$$
(9)

Since Δs is uniform for interior nodes of the grid, substitutions like the following have been made for simplicity's sake:

$$k_{\mathbf{NP}} = \frac{2k_{\mathbf{N}}k_{\mathbf{P}}}{k_{\mathbf{N}} + k_{\mathbf{P}}}.$$
 (10)

Equation 8 describes the temperature change for one node P over one time step. The values ρ , $C_{\mathbf{p}}$ and k are held constant for each time step, but may be changed between steps. It is necessary to apply eq 8 (or the appropriate boundary equation) to each node of the grid to solve for the temperature distribution for one time step. The process is repeated until the desired number of time steps have been completed. Let us now investigate several methods used in the solution of the set of equations.

The first and probably the easiest method to program is the explicit method, where eq 8 is solved for T_P' (which is the value of T_P for time $t + \Delta t$). Note that all of the other temperature variables represent time t, so that the equation may be solved explicitly. An initial temperature is assigned to each node in the grid; the appropriate equation is solved for each node to determine the temperature resulting from the first time step. The resultant temperature distribution is then used in calculations for the second time step. This process continues until the desired number of time steps have been calculated. The main problem with this method is that a stability criterion must be met to produce a reliable answer. For a two-dimensional homogeneous grid, the criterion is

$$\frac{(\Delta s)^2 \rho C_p}{k\Delta t} > 4.$$

The second secon

Note that a balance must be struck between the internodal distance (Δs) and the time step (Δt). In problems involving steep temperature gradients, a small internodal spacing must be used; consequently, the time step must be very small. The combination of many nodes and time steps makes the solution costly in computer time.

A second method use J in solving finite difference equations is the implicit method. Here T_P' is taken as the value of T_P for time t, and all of the other variables represent time $t + \Delta t$; hence, all temperature variables are unknown except T_P' . An appropriate equation must be written for each node in the grid, and the resulting set of equations is solved simultaneously for each time step. For the two-dimensional case of an x by y grid, there will be xy equations in xy unknowns. The matrix of coefficients is banded (bandwidth = 2x + 1), but in general there is no symmetry about the main diagonal. Large matrices often occur and they are usually solved by an iterative method. The implicit method has the advantage of being unconditionally stable, but it may require a fair number of iterations for adequate convergence in the matrix solution.

The finite difference method chosen for this computer program is known as the alternating direction implicit method (Peaceman and Rachford 1955, Carnahan et al. 1969). It is unconditionally stable like the fully implicit method, yet it requires only the solution of a tridiagonal matrix, thus being efficient in computer time and computer core storage. (The matrix-solving algorithm is discussed in the TRIDIG, The Matrix Solver section). Essentially, the method requires the solution of two different equations for each time step. The first equation is implicit only in the horizontal direction. The results from the first step are then used to solve the second equation, which is implicit only in the vertical direction.

For example, apply the alternating direction considerations to a node in the interior of the grid. Equation 9, which describes the heat balance for an interior node, is repeated here for convenience:

$$k_{\rm NP} (T_{\rm N} - T_{\rm P}) + k_{\rm WP} (T_{\rm W} - T_{\rm P}) + k_{\rm SP} (T_{\rm S} - T_{\rm P}) + k_{\rm EP} (T_{\rm E} - T_{\rm P}) = \frac{(\Delta s)^2 \rho' C_{\rm P}'}{\Delta t} T_{\rm p}' - \frac{(\Delta s)^2 \rho C_{\rm P}}{\Delta t} T_{\rm p}.$$

Each node in the grid is given an initial temperature. Then, for the first pass, let the above equation be implicit in the horizontal direction. The time step for the pass will be $\Delta t/2$, and the above equation may be written as follows:

$$k_{WP}T'_{W} + k_{EP}T'_{E} - \left[k_{WP} + k_{EP} + \frac{2(\Delta s)^{2} \rho' C'_{P}}{\Delta t}\right] T'_{P} = -k_{NP}T_{N} - k_{SP}T_{S} + \left[k_{NP} + k_{SP} - \frac{2(\Delta s)^{2} \rho C_{P}}{\Delta t}\right] T_{P}.$$
(11)

The primed variables represent the value of those variables at time $t - \Delta t/2$. The right-hand side of the equation is known; T'_W , T'_P and T'_E are to be determined. Equation 11 (or the appropriate boundary equation) is applied to each node of a row of the grid, creating a tridiagonal matrix of coefficients. Such a system of equations is quickly solved for each row of the grid. The resulting temperature distribution represents $t + \Delta t/2$, the end of the first pass.

The process is similar for the second pass, except that now eq 9 is rewritten to be implicit in the vertical direction:

$$k_{NP}T'_{N} + k_{SP}T'_{S} - \left[k_{NP} + k_{SP} + \frac{2(\Delta s)^{2} \rho' C'_{P}}{\Delta t}\right] T'_{P} = -k_{WP}T_{W} - k_{EP}T_{E} + \left[k_{WP} + k_{EP} - \frac{2(\Delta s)^{2} \rho C_{P}}{\Delta t}\right] T_{P}.$$
(12)

The primed variables represent time $t + \Delta t$, and the others represent time $t + \Delta t/2$. The tridiagonal matrix is formed and solved for each column of the grid; the resultant temperatures represent the temperature distribution in the grid after $t + \Delta t$, one time step. The entire process is repeated until the desired number of steps have been calculated.

The reader may consult Carnahan et al. (1969), Holman (1972), Croft and Lilley (1977) and Mitchell and Griffiths (1980) for more information on finite difference methods.

Boundary conditions

The following boundary conditions will be derived from heat balance considerations. In this paper each boundary equation will not be expressed in the form needed for use in the alternating direction procedure, but the reader should be aware that each equation was put into that form for use in the computer programs ADI and ADIPC.

Sides of the grid

Consider a node, P, on the right-hand grid boundary. The control volume associated with the node is the area enclosed by dotted lines around it, as depicted in Figure 3.

The sum of the heat flowing through the boundaries of P's control volume is responsible for the temperature change of the node, that is,

$$Q_{\rm NP} + Q_{\rm WP} + Q_{\rm SP} + Q_{\rm SP} = \frac{\Delta x \, \Delta y \, \rho \, C_{\rm P} \, \Delta T_{\rm P}}{2\Delta t} \tag{13}$$

where Q_{NP} = the heat flow from node N to (or from) node P

 $\frac{\Delta x}{2} \cdot \Delta y$ = the volume (for unit depth) of node P

 ρ = the density of the material in node P

 $C_{\rm p}$ = the specific heat

T = temperature

t = time.

The flow of heat from node N to node P is given by

$$Q_{NP} = \frac{1}{2} k_{NP} (T_N - T_P)$$
 (14)

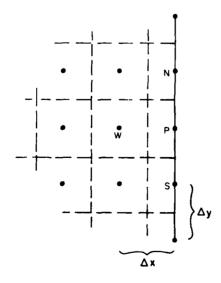


Figure 3. A node on a right-hand boundary.

where k_{NP} is the effective conductivity between nodes N and P, and for $\Delta x = \Delta y = \Delta s$,

$$k_{\rm NP} = \frac{2k_{\rm N}k_{\rm P}}{(k_{\rm N} + k_{\rm P})}$$

Similarly, the heat flow from nodes W and S to node P may be given as

$$Q_{\mathbf{WP}} = k_{\mathbf{WP}} \left(T_{\mathbf{W}} - T_{\mathbf{P}} \right) \tag{15}$$

$$Q_{\rm SP} = \frac{1}{2} k_{\rm SP} (T_{\rm S} - T_{\rm P}). \tag{16}$$

Now consider several cases describing the heat flow across the boundary.

Constant flux boundary. For a boundary subject to a constant heat flux, the condition

$$\frac{\partial T}{\partial x}\Big|_{x=P} = \text{constant}$$

exists at the boundary. For node P, illustrated in Figure 3, the heat flow across its east side is given by

$$Q_{\rm EP} = \phi \cdot \Delta y \tag{17}$$

where ϕ is the heat flux per unit area crossing the boundary. The equation for a node on the right-hand side of the grid with constant flux is obtained by combining eq 13-17 and allowing $\Delta x = \Delta y = \Delta s$:

$$\frac{1}{2} k_{NP} (T_N - T_P) + k_{WP} (T_W - T_P) + \frac{1}{2} k_{SP} (T_S - T_P) + \phi \Delta s = \frac{(\Delta s)^2}{2\Delta t} [\rho' C_P' T_P' - \rho C_P T_P].$$
(18)

The indices may be suitably rearranged for constant flux boundaries on other sides of the grid.

Note that for a boundary which is insulated or on a line of symmetry the zero heat flux condition holds, and $\phi = 0$.

Convection boundary. For the node illustrated in Figure 3, exposed to convection on the right-hand side, the heat flow from the convective medium to node P is given by

$$Q_{\rm EP} = h_{\rm E} \Delta y \left(T_{\rm E} - T_{\rm P} \right) \tag{19}$$

where $h_{\rm E}$ is the coefficient of convective heat transfer and $T_{\rm E}$ is the temperature outside the grid. Combining equations 13, 14, 15, 16 and 19, we arrive at the equation for a node on the convective right-hand side boundary, for $\Delta x = \Delta y = \Delta s$:

$$\frac{1}{2} k_{NP} (T_N - T_P) + k_{WP} (T_W - T_P) + \frac{1}{2} k_{SP} (T_S - T_P) + h\Delta s (T_E - T_P) = \frac{(\Delta s)^2}{2\Delta t} [\rho' C_P' T_P' - \rho C_P T_P].$$
(20)

Specified temperature boundary. For a node of specified temperature on a boundary, corner or inside the grid, apply the equation $T_p = C$, where C is the temperature of the node at time t.

Semi-infinite boundary. This condition represents a continuous, uniform material extending in one direction, with a known temperature a large distance away. It is approximated here by use

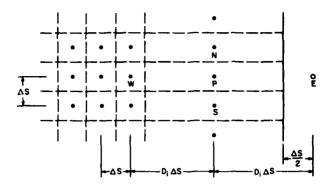


Figure 4. A node on a semi-infinite boundary.

of a large internodal distance between the last two nodes of a row of the grid. Consider the situation illustrated in Figure 4, for the right-hand side of the grid. Node E is not actually a part of the grid, but is the location of known temperature $T_{\rm E}$ outside the grid.

The distance between the last two nodes and also between the boundary node and the location of the known temperature T_E is $D_i \Delta s$. Note that D_i represents a multiple of Δs . The heat flow from nodes N, W, S and E to node P may be calculated as follows

$$Q_{NP} = k_{NP} (2D_i - 1) (T_N - T_P)$$
 (21)

$$Q_{\mathbf{WP}} = k_{\mathbf{WP}} \left(T_{\mathbf{W}} - T_{\mathbf{P}} \right) \tag{22}$$

$$Q_{\rm SP} = k_{\rm SP} (2D_{\rm i} - 1) (T_{\rm S} - T_{\rm P}) \tag{23}$$

$$Q_{\rm EP} = k_{\rm EP} \left(T_{\rm E} - T_{\rm P} \right). \tag{24}$$

The effective conductivities are of the same form as eq 10, except for k_{WP} and k_{EP}

$$k_{WP} = \frac{2k_W k_P}{[k_P + (2D_i - 1) k_W]}$$
.

Node E is assumed to be the same material as node P, thus $k_{\rm EP} = k_{\rm P}$. The sum of the heat flows accounts for the temperature change of the node, thus the equation for a node on the right-hand semi-infinite boundary is

$$k_{NP}(2D_i - 1)(T_N - T_P) + k_{WP}(T_W - T_P) + k_{SP}(2D_i - 1)(T_S - T_P) +$$

$$k_{\rm EP} (T_{\rm E} - T_{\rm P}) = \frac{(2D_{\rm i} - 1) (\Delta s)^2}{\Delta t} (\rho' C_{\rm P}' T_{\rm P}' - \rho C_{\rm P} T_{\rm P})$$
 (25)

where $(2D_i - 1)(\Delta s)^2$ is the control volume for the node per unit depth.

In any finite difference formulation, the accuracy of the solution increases as the area represented by a node decreases. Therefore, when using the semi-infinite boundary formulation, the user should specify the smallest D_i acceptable. The semi-infinite condition should only be used in regions where the temperature gradient is small and precise knowledge of the temperature distribution is not critical.

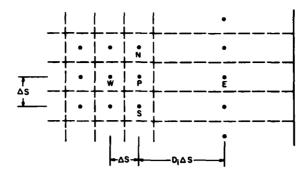


Figure 5. An interior node adjacent to the semi-infinite boundary.

This boundary conditions also requires a special heat balance for the node adjacent to the semi-infinite node. Consider again the situation for a right-hand boundary, as illustrated in Figure 5. The appropriate heat flow equations are

$$Q_{NP} = k_{NP} \left(T_N - T_P \right) \tag{26}$$

$$Q_{\mathbf{WP}} = k_{\mathbf{WP}} \left(T_{\mathbf{W}} - T_{\mathbf{P}} \right) \tag{27}$$

$$Q_{\rm SP} = k_{\rm SP} \left(T_{\rm S} - T_{\rm P} \right) \tag{28}$$

where k_{NP} , k_{WP} , k_{SP} are of the same form as eq 10, and

$$Q_{\rm EP} = k_{\rm EP} \left(T_{\rm N} - T_{\rm P} \right) \tag{29}$$

where $k_{\rm EP} = \frac{2k_{\rm E}k_{\rm P}}{\left[k_{\rm E} + (2D_{\rm i} - 1)\,k_{\rm P}\right]}$.

The equation for an interior node adjacent to a right-hand semi-infinite boundary node is

$$k_{NP} (T_{N} - T_{P}) + k_{WP} (T_{W} - T_{P}) + k_{SP} (T_{S} - T_{P}) + k_{EP} (T_{E} - T_{P}) = \frac{(\Delta s)^{2}}{\Delta t} [\rho' C_{P}' T_{P}' - \rho C_{P} T_{P}].$$
(30)

Corner of the grid

The corners of the grid require special heat balances, depending on the conditions on each edge of the grid. Consider the upper right-hand corner of the grid, shown in Figure 6. $T_{\rm N}$ and $T_{\rm E}$ are not a part of the grid, but are temperatures outside the grid. The heat flow from the two nodes adjacent to node P may be given as

$$Q_{WP} = \frac{1}{2} k_{WP} (T_W - T_P)$$
 (31)

$$Q_{\rm SP} = \frac{1}{2} k_{\rm SP} (T_{\rm S} - T_{\rm P}) \tag{32}$$

where k_{WP} and k_{SP} follow the form given in eq 10. Q_{NP} and Q_{EP} are dependent upon the

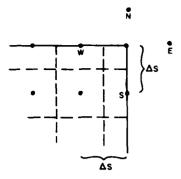


Figure 6. A node on a corner of the grid.

particular boundary conditions and will be presented shortly. The corner equations will follow the form

$$Q_{\rm NP} + Q_{\rm WP} + Q_{\rm SP} + Q_{\rm EP} = \frac{(\Delta s)^2}{4} \rho C_{\rm P} \frac{\Delta T}{\Delta t}$$
 (33)

where $\frac{(\Delta s)^2}{4}$ is the volume represented by node P for unit depth.

Constant flux on both sides. Let ϕ_N be the heat flux per unit area crossing the north side of the corner shown in Figure 6, and ϕ_E be that crossing the east side of the corner. Then the heat flows across the two sides are given by

$$Q_{\rm NP} = \phi_{\rm N} \cdot \frac{\Delta s}{2} \tag{34}$$

$$Q_{\rm EP} = \phi_{\rm E} \cdot \frac{\Delta s}{2} \ . \tag{35}$$

The equation for the corner is

$$\frac{1}{2} k_{WP} (T_W - T_P) + \frac{1}{2} k_{SP} (T_S - T_P) + \frac{1}{2} \Delta s (\phi_N + \phi_E) = \frac{(\Delta s)^2}{4\Delta t} [\rho' C_P' T_P' - \rho C_P T_P].$$
(36)

As in the case of the side of the grid, if a side of the corner is insulated or on a line of symmetry, $\phi = 0$ for that side.

Convection on both sides. For a corner subject to convection on both of its sides, the heat flow from each of the two sides may be given as

$$Q_{NP} = h_N \frac{\Delta s}{2} \left(T_N - T_P \right) \tag{37}$$

$$Q_{\rm EP} = h_{\rm E} \frac{\Delta s}{2} \quad (T_{\rm E} - T_{\rm P}) \tag{38}$$

where h_N is the coefficient of convective heat transfer on the north side and h_E is that on the east side. The equation for node P is

$$\frac{1}{2} k_{WP} (T_W - T_P) + \frac{1}{2} k_{SP} (T_S - T_P) + \frac{1}{2} h_N \Delta s (T_N - T_P) + \frac{1}{2} h_E \Delta s (T_E - T_P) =$$

$$\frac{(\Delta s)^2}{4\Delta t} \left[\rho' C_P' T_P' - \rho C_P T_P \right].$$
(39)

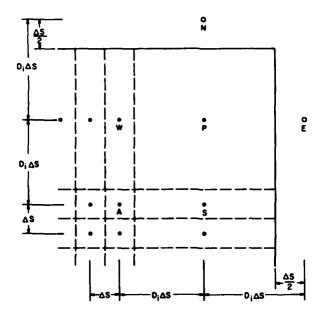


Figure 7. A semi-infinite corner.

Semi-infinite on both sides. For a corner with semi-infinite conditions on both sides, an irregular finite difference grid is employed once more. Consider again the upper right-hand corner, now illustrated in Figure 7. $T_{\rm N}$ and $T_{\rm E}$ are known temperatures outside the grid. Heat flows are given by

$$Q_{NP} = k_{NP} (2D_i - 1) (T_N - T_P)$$
(40)

$$Q_{WP} = k_{WP} (2D_i - 1) (T_W - T_P)$$
(41)

$$Q_{\rm SP} = k_{\rm SP} (2D_{\rm i} - 1) (T_{\rm S} - T_{\rm P}) \tag{42}$$

$$Q_{EP} = k_{EP} (2D_i - 1) (T_E - T_P).$$
(43)

Again,

$$k_{\text{NP}} = k_{\text{EP}} = k_{\text{P}}; k_{\text{WP}} = \frac{2k_{\text{W}}k_{\text{P}}}{k_{\text{P}} + (2D_i - 1)k_{\text{W}}} \text{ and } k_{\text{SP}} = \frac{2k_{\text{S}}k_{\text{P}}}{k_{\text{P}} + (2D_i - 1)k_{\text{S}}}.$$

Then the equation for a semi-infinite corner node follows:

$$k_{\rm NP} \, (2D_{\rm i} - 1) \, (T_{\rm N} - T_{\rm P}) + k_{\rm WP} \, (2D_{\rm i} - 1) \, (T_{\rm W} - T_{\rm P}) + k_{\rm SP} \, (2D_{\rm i} - 1) \, (T_{\rm S} - T_{\rm P}) + k_{\rm SP} \, (2D_{\rm i} - 1) \, (T_{\rm N} - T_{\rm P}) + k_{\rm SP} \, (2D_{\rm i} - 1) \, (T_{\rm N} - T_{\rm P}) + k_{\rm NP} \, (2D_{\rm i} - 1) \, (T_{\rm N} - T_{\rm$$

$$k_{\rm EP} (2D_{\rm i} - 1) (T_{\rm E} - T_{\rm P}) = \frac{[(2D_{\rm i} - 1) (\Delta s)]^2}{\Delta t} [\rho' C_{\rm P}' T_{\rm P}' - \rho C_{\rm P} T_{\rm P}].$$
 (44)

Because the nodes are of irregular size, special consideration must be given also to node temperatures $T_{\rm W}$, $T_{\rm S}$ and $T_{\rm A}$ in Figure 7.

Consider first node A. This is similar to the node P illustrated in Figure 5, except that the distance to node N is now $D_i \cdot \Delta s$. This results in our changing k_{NP} to

$$k_{\rm NP} = \frac{2k_{\rm N}k_{\rm P}}{k_{\rm N} + (2D_{\rm i} - 1)k_{\rm P}}$$

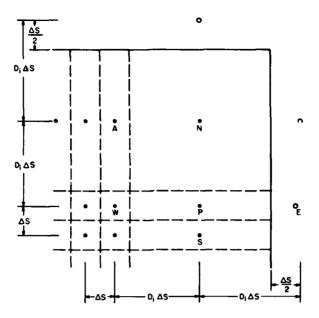


Figure 8. A semi-infinite side node adjacent to a semi-infinite corner.

Except for this change in k_{NP} , the equation for the square interior node inside an upper right-hand corner with semi-infinite conditions on both sides is the same as eq 29.

Now consider the semi-infinite nodes adjacent to the semi-infinite corner node. These are node temperatures $T_{\mathbf{W}}$ and $T_{\mathbf{S}}$ in Figure 7. The figure is presented again, this time as Figure 8, with the nodes relabeled. The type of node under consideration is represented by nodes A and P in Figure 8. The following equations reflect the heat flow to node P:

$$Q_{NP} = k_{NP} (2D_i - 1) (T_N - T_P)$$
(45)

$$Q_{\mathbf{WP}} = k_{\mathbf{WP}} \left(T_{\mathbf{W}} - T_{\mathbf{P}} \right) \tag{46}$$

$$Q_{SP} = k_{SP} (2D_i - 1) (T_S - T_P)$$
(47)

$$Q_{\rm FP} = k_{\rm FP} \left(T_{\rm S} - T_{\rm P} \right). \tag{48}$$

Here,

$$k_{\text{NP}} = \frac{2k_{\text{N}}k_{\text{P}}}{k_{\text{N}} + (2D_{\text{i}} - 1)k_{\text{P}}}$$
 and $k_{\text{WP}} = \frac{2k_{\text{W}}k_{\text{P}}}{k_{\text{P}} + (2D_{\text{i}} - 1)k_{\text{W}}}$

The equation for this node is

$$k_{NP}(2D_i-1)(T_N-T_P)+k_{WP}(T_W-T_P)+k_{SP}(2D_i-1)(T_S-T_P)+$$

$$k_{\rm EP} (T_{\rm E} - T_{\rm P}) = \frac{(2D_{\rm i} - 1) (\Delta s)^2}{\Delta t} [\rho' C_{\rm P}' T_{\rm P}' - \rho C_{\rm P} T_{\rm P}]. \tag{49}$$

Constant flux, convective corner. For the corner illustrated in Figure 6, allow convection to occur across the north side of node P and constant flux across the east side

$$Q_{NP} = h_N \frac{\Delta s}{2} \left(T_N - T_P \right) \tag{50}$$

 $Q_{\rm EP}$ is given by eq 35, and $Q_{\rm WP}$ and $Q_{\rm SP}$ are given by eq 31 and 32, respectively. The resulting equation for the corner is:

$$\frac{1}{2} h_{N} \Delta s (T_{N} - T_{P}) + \frac{1}{2} k_{WP} (T_{W} - T_{P}) + \frac{1}{2} k_{SP} (T_{S} - T_{P}) + \frac{1}{2} \phi_{E} \Delta s = \frac{(\Delta s)^{2}}{4 \Delta t} [\rho' C_{P}' T_{P}' - \rho C_{P} T_{P}].$$
(51)

Constant flux, semi-infinite corner. The corner illustrated in Figure 9 has a semi-infinite boundary on the right side. Let there be a constant flux per unit area, ϕ_N , across the north side of node P, then

$$Q_{NP} = \phi_N (2D_i - 1) \Delta s \tag{52}$$

$$Q_{WP} = \frac{1}{2} k_{WP} (T_W - T_P)$$
 (53)

$$Q_{\rm SP} = k_{\rm SP} (2D_{\rm i} - 1) (T_{\rm S} - T_{\rm P}) \tag{54}$$

$$Q_{\rm EP} = \frac{1}{2} k_{\rm EP} (T_{\rm E} - T_{\rm P})$$

$$k_{WP} = \frac{2k_W k_P}{k_P + (2D_i - 1) k_W}$$

$$k_{\rm SP} = \frac{2k_{\rm S}k_{\rm P}}{k_{\rm S} + k_{\rm P}}$$

$$k_{\rm EP} = k_{\rm P} \ . \tag{55}$$

The equation for the corner is as follows:

$$\phi_{N} (2D_{i}-1) \Delta s + \frac{k_{WP}}{2} (T_{W}-T_{P}) + k_{SP} (2D_{i}-1) (T_{S}-T_{P}) + \frac{k_{EP}}{2} (T_{E}-T_{P}) = \frac{(2D_{i}-1) (\Delta s)^{2}}{2\Delta t} [\rho' C_{P}' T_{P}' - \rho C_{P} T_{P}].$$
(56)

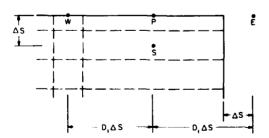


Figure 9. A node on a semi-infinite corner.

Again, the node adjacent to the semi-infinite node must have a special equation. For node W in Figure 9, with constant flux ϕ_N across the top, eq 57 applies, with reference to the node under consideration as node P:

$$\phi_{N} \Delta s + \frac{1}{2} k_{WP} (T_{W} - T_{P}) + k_{SP} (T_{S} - T_{P}) + \frac{1}{2} k_{EP} (T_{E} - T_{P}) = \frac{(\Delta s)^{2}}{2\Delta t} \left[\rho' C_{P}' T_{P}' - \rho C_{P} T_{P} \right].$$
(57)

Semi-infinite, convective corner. For the corner in Figure 9, allow convection across the north side, and allow the same semi-infinite boundary on the east side,

$$Q_{NP} = h_N (2D_i - 1) (T_N - T_P). (58)$$

 Q_{WP} , Q_{SP} and Q_{EP} appear in eq 53, 54 and 55, respectively. Thus the equation for the configuration is

$$h_{N} (2D_{i}-1) (T_{N}-T_{P}) + \frac{k_{WP}}{2} (T_{W}-T_{P}) + k_{SP} (2D_{i}-1) (T_{S}-T_{P}) + \frac{k_{EP}}{2} (T_{E}-T_{F}) = \frac{(2D_{i}-1) (\Delta s)^{2}}{2\Delta t} [\rho' C_{P}' T_{P}' - \rho C_{P} T_{P}].$$
(59)

The appropriate equation for node W in Figure 9, subject to convection across the top, follows. Again, refer to this node as node P:

$$h_{N} \Delta s (T_{N} - T_{P}) + \frac{1}{2} k_{WP} (T_{W} - T_{P}) + k_{SP} (T_{S} - T_{P}) + \frac{k_{WP}}{2} (T_{W} - T_{P}) = \frac{(\Delta s)^{2}}{2\Delta t} [\rho' C_{P}' T_{P}' - \rho C_{P} T_{P}].$$
(60)

PHASE CHANGE

The methods discussed in this report so far apply to both ADIPC and ADI, the programs developed for heat conduction with and without phase change. The ideas on phase change presented in this section will apply only to ADIPC. The computer program solves the heat conduction equation only; possible effects of moisture migration through the medium (and unfrozen water content in a frozen soil, for example) are neglected. Although the examples used involve freezing, ADIPC permits either freezing or thawing at any node of the grid.

During the freezing process, the temperature of a substance that is initially above its fusion temperature decreases as heat is removed until the substance reaches its fusion temperature. The continuing removal of heat produces no change in temperature until the energy equivalent to the latent heat of the substance is removed. The substance has then changed state to become frozen; any further removal of heat again results in a decrease of temperature below the temperature of fusion.

ADIPC models freezing or thawing by defining an apparent specific heat during phase change that accounts for the entire enthalpy change that takes place, including the enthalpy in the latent heat of fusion. To do this, it must be assumed that phase change takes place over a finite temperature range, ΔT , around the fusion temperature. Allow T_f to denote the fusion temperature, C_{P_A}

and C_{P_B} to be the specific heat of the substance above and below freezing, respectively, and let C_{P}^{*} be the apparent specific heat defined for phase change. For a temperature change in a substance not involving a change of state, the enthalpy change ΔH is given as

$$\Delta H = \int C_{\mathbf{P}} dT.$$

Now the temperature range for phase change will be $T_f \pm \Delta T/2$, and the apparent specific heat accounts for the entire enthalpy change as follows, where H_L represents the latent heat of fusion:

$$\int_{T_{\rm f}}^{T_{\rm f}} + \frac{\Delta T}{2} C_{\rm p}^{*} dT = \int_{T_{\rm f}}^{T_{\rm f}} C_{\rm p_B} dT + \int_{T_{\rm f}}^{T_{\rm f}} + \frac{\Delta T}{2} C_{\rm p_A} dT + H_{\rm L}.$$
 (61)

 ΔT is chosen small enough so that each specific heat may be assumed constant over the integral,

$$C_{\mathbf{p}}^{*} \Delta T = C_{\mathbf{p}_{\mathbf{B}}} \frac{\Delta T}{2} + C_{\mathbf{p}_{\mathbf{A}}} \frac{\Delta T}{2} + H_{\mathbf{L}}$$

$$C_{\mathbf{p}}^{*} = \frac{1}{2} \left(C_{\mathbf{p}_{\mathbf{B}}} + C_{\mathbf{p}_{\mathbf{A}}} \right) + \frac{H_{\mathbf{L}}}{\Delta T} . \tag{62}$$

Equation 62 defines the apparent heat capacity as used in ADIPC. The apparent heat capacity method is further discussed in Bonacina and Comini (1973).

Phase change is implemented in ADIPC as follows. At the beginning of each time step, each node of the grid is examined. If its temperature lies within the range $T_f \pm \Delta T/2$, the specific heat for that node is defined by eq 62. The conduction equations are set up and solved as usual. At the end of the complete time step, the newly calculated temperature of each node is compared with the temperature of that node at the beginning of the time step. If the temperature skipped from below $T_f - \Delta T/2$ to above $T_f + \Delta T/2$, or vise versa, then the phase change front skipped that node, and the temperature of the node is reassigned as follows. If the temperature of the node skipped from the frozen to the unfrozen domain,

$$T = \left(T_{\rm f} - \frac{\Delta T}{2}\right) + \frac{C_{\rm P_B}}{C_{\rm p}^{\bullet}} \quad \left[T' - \left(T_{\rm f} - \frac{\Delta T}{2}\right)\right] \tag{63}$$

where T' is the calculated nodal temperature before its reassignment. If the node skipped from an unfrozen state to the frozen state, its temperature is reassigned as follows:

$$T = \left(T_{\rm f} + \frac{\Delta T}{2}\right) + \frac{C_{\rm P_A}}{C_{\rm P}^*} \left[T' - \left(T_{\rm f} + \frac{\Delta T}{2}\right)\right]. \tag{64}$$

In this way, the program assures that the phase change front does not skip a node.

A drawback with the apparent heat capacity method is that it is not designed to follow the exact location of the phase change front for each time step; it is designed to calculate nodal temperature. The location of the $T_{\rm f}$ isotherm may be found by interpolation, but it approximates the location of the front in a step-like pattern because of the discretization of a continuous space, as will be illustrated in the program verifications. This behavior is minimized by use of a small internodal spacing. The method is flexible enough to deal with phase change in two-dimensional space when there may be several locations of phase change fronts; for this flexibility the method was chosen.

COMPUTER PROGRAM

In using finite differences, we replace differentials and derivatives with differences. It makes sense, then, that the accuracy of the solution increases as these differences become smaller. Specifically, the accuracy improves as the nodal spacing and time step, Δs and Δt , are made small. There is no set criterion for just how small they must be, usually this is discovered through trial and error.

The grid for these programs may be any size, so far as the outside dimensions are concerned, but, except for the semi-infinite boundary conditions, each node in the grid represents a square area of unit depth. RAY (I, J, K) is the array that represents the grid. I and J are spatial variables in Cartesian coordinates; I is incremented vertically and J horizontally. K has three values. RAY (I, J, 1) represents the temperature distribution in the grid, i.e., a temperature is stored at each I, J location. RAY (I, J, 2) records the nodal location type. Examples of location types include a node on a constant flux boundary, variable interior node, specified temperature corner node, etc. Each location type is assigned a number, and one such number is stored for each I, J location under RAY (I, J, 2). This information is used to assign the correct form of the heat conduction equation to each node. Each material type in the problem is given a number; these numbers are stored in RAY (I, J, 3) and may be used to assign a conductivity, density and specific heat for each node.

The programs are made up of four parts: 1) a data-gathering subroutine, 2) the main program, 3) a subroutine to solve the tridiagonal matrix and 4) a subroutine to locate the user-specified isotherms at times specified by the user. Parts 3 and 4 are identical for both ADI and ADIPC, and part 1 is similar for each. Therefore, parts 1, 3 and 4 will be discussed only once but are included in each of the two programs.

ADDATA, the data subroutine

The necessary data for the program are handled mainly through subroutine ADDATA. Variables used are defined in the comment statements at the beginning of the subroutine. The user has simply to edit the subroutine, following directions in the comment statements in the subroutine to initialize variables and arrays. When the subroutine is run, the data are put into a formatted data file ADIDAT; the user does not have to worry about formats.

TRIDIG, the matrix solver

This subroutine solves the tridiagonal matrix formed in the main program. The matrix is formed by the application of eq 11 or 12 (or a boundary condition counterpart) to each node of the grid. The resultant system of equations is illustrated in matrix form in Figure 10. The matrix of coefficients is an n by n matrix; all entries are zero except those on the three center diagonals, hence the term tridiagonal matrix. The vector on the right contains the elements of the right-hand side of each equation. To conserve computer storage space, only the three diagonals are stored, as vectors. Thus the effective storage space required is reduced from n by n + 1 to n by 4.

The solution algorithm is commonly used (e.g. Gerald 1980). Each element of the lower diagonal may be eliminated by subtracting the appropriate multiple of the (i-1) row from the *i*th row. The values of b_i and d_i , after elimination of a_i , are

$$b_{i} = b_{i} - \left(\frac{a_{i}}{b_{i-1}}\right) c_{i-1} \tag{65}$$

$$d_i = d_i - \left(\frac{a_i}{b_{i-1}}\right) d_{i-1}$$
 for $i = 1, 2, 3 \dots n$. (66)

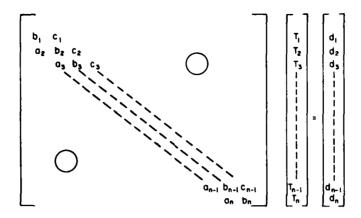


Figure 10. Tridiagonal matrix equation,

After replacing the element of vectors b and d by the new values, we perform a back substitution as follows:

$$d_{n} = \frac{d_{n}}{b_{n}} \tag{67}$$

$$d_{i} = \frac{d_{i} - c_{i}d_{i+1}}{b_{i}}, \quad \text{for } i = n - 1, n - 2, \dots 1.$$
 (68)

The elements of the solution vector replace vector d.

ISOTHM, the isotherm finder

Subroutine ISOTHM examines the temperature distribution in the grid and performs a linear interpolation between adjacent nodes to produce the Cartesian coordinates of the locations of the user-specified isotherms. The coordinates are listed in file POINTS in two columns, one for the horizontal coordinate and one for the vertical coordinate. POINTS may be used as a data file for a plotting routine. Subroutine ISOTHM may be called at the completion of any number of time steps; the frequency is specified by the user in ADDATA.

ADI, main program

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ADI first interactively asks the user if ADDATA should be called. The subroutine should be called if ADIDAT does not already contain the formatted data. If the user so desires ADDATA is called. If not, ADI reads the data from ADIDAT.

For each time step, the thermal properties such as conductivity, density and specific heat are updated, if appropriate, and the resultant conductivities between the nodes are figured. Each node is assigned the coefficients for the appropriate form of the conduction equation, forming the tridiagonal matrix (Fig. 10), and subroutine TRIDIG is called to solve the matrix. The resultant temperature distribution is then used in the second half-time step, and TRIDIG is called once again. These temperatures then represent the distribution for one complete time step. If appropriate, subroutine ISOTHM is then called to locate the user-specified isotherms. This procedure is repeated for the desired number of time steps.

The initial data and boundary conditions are printed into file ADIOUT; the temperature distributions for specified time steps are printed into file ADITMP, and finally, a new data file, ADNDAT, is created for the final temperature distribution. This file is in the same format as the input file ADIDAT; ADNDAT may be used as a starting point if the user wishes to run the model for more time.

ADIPC, main program

ADIPC is similar to ADI, but differs in that phase change considerations are implemented. At the beginning of each time step, if the temperature of a node lies in the phase change temperature range, the program assigns the apparent specific heat to that node, as previously discussed. Otherwise, the specific heat, thermal conductivity and density for each node are updated according to the user's specifications.

At the end of the time step, the temperature of each node is compared to its temperature at the end of the previous time step. If the temperature skipped over the phase change temperature range, it is reassigned as discussed in the *Phase Change* section. Then, if appropriate, subroutine ISOTHM is called.

The initial and boundary conditions are printed in readable form in file ADPOUT; the temperature distributions for specified time steps are printed into file ADPTMP, and a new data file with the final temperature distributions is printed into file ADPNDT.

VERIFICATION OF ADI

Comparison of ADI with analytical results

Semi-infinite corner

The results of ADI will first be compared to the problem of a semi-infinite corner, as illustrated in Figure 11.

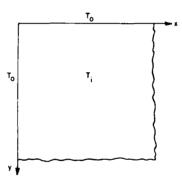


Figure 11. Semi-infinite corner problem.

 T_1 is the uniform initial temperature; at time t = 0 the temperature of the two edges, given by x = 0 and y = 0, is changed to T_0 . The solution is well documented (Carslaw and Jaeger 1959, and Holman 1972) and is found by using a product solution for two one-dimensional problems. The solution is given by

$$\frac{T - T_0}{T_1 - T_0} = \operatorname{erf} \frac{x}{2\sqrt{\alpha t}} - \operatorname{erf} \frac{y}{2\sqrt{\alpha t}}$$
(69)

where α is the thermal diffusivity, and

$$\operatorname{erf} \frac{x}{2\sqrt{\alpha t}} = \frac{2}{\sqrt{\pi}} \int_{0}^{x/(2\sqrt{\alpha t})} e^{-\eta^2} d\eta.$$

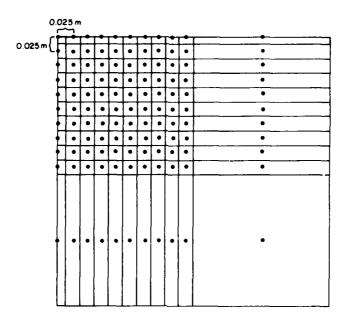


Figure 12. Finite difference grid for the semi-infinite corner problem.

For a comparison run, the following values were used:

$$T_i = 20^{\circ} \text{C (68°F)}$$

 $T_0 = 40^{\circ} \text{C (104°F)}$
 $\alpha = 0.0025 \text{ m}^2/\text{hr}$

The finite difference grid appears as shown in Figure 12. The boundary condition assigned to the top and left side was that of constant temperature equal to 40° C. The right side and bottom of the grid were assigned a semi-infinite boundary condition, with $T_{\rm S} = T_{\rm E} = 20^{\circ}$ C and $D_{\rm i} = 50$ (see the Sides of the Grid section). The nodes not situated on a boundary were assigned an initial temperature of 20° C. A time step of 0.25 hr was used, and the internodal spacing was 0.025 m.

The locations of the 35°, 30° and 25°C isotherms are plotted for the results of ADI and for the analytical solution for several time steps in Figure 13. Excellent agreement is found between the two solutions for regions not adjacent to the semi-infinite boundary. As previously stated, the semi-infinite condition is an approximation. Also, the large internodal distance used for the semi-infinite boundary decreases the accuracy of the solution in that region.

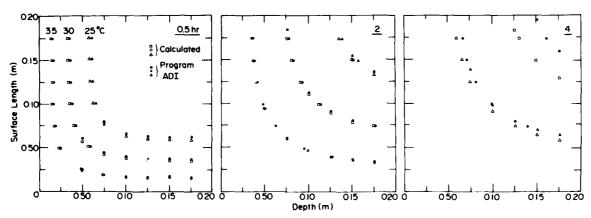


Figure 13. Locations of 35°, 30° and 25° C isotherms from ADI and from the analytical solution for the semi-infinite corner problem.

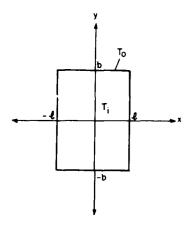


Figure 14. Finite rectangle problem.

Finite rectangle

In order to provide a two-dimensional comparison that doesn't involve the use of semi-infinite boundaries, and also to demonstrate internodal spacing effects, consider the problem of a rectangle of uniform initial temperature T_i that is subject to a step change in the temperatures of all of the edges to T_0 at time zero. The problem is illustrated in Figure 14. The solution of this problem is again obtainable from the product solution of two one-dimensional problems. Carslaw and Jaeger (1959) provide the one dimensional solution:

$$\frac{T(x,t)-T_0}{T_i-T_0} = \frac{4}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{2n+1} \exp \frac{-\alpha (2n+1)^2 \pi^2 t}{4\ell^2} \cos \frac{(2n+1) \pi x}{2\ell} . \tag{70}$$

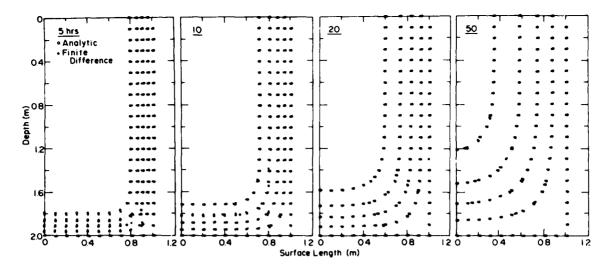
The product solution is then given by

$$T(x,y) = T_0 + \frac{16(T_1 - T_0)}{\pi^2} \sum_{n=0}^{\infty} \frac{(-1)^n}{2n+1} \exp \frac{-\alpha (2n+1)^2 \pi^2 t}{4\ell^2} \cos \frac{(2n+1)\pi x}{2\ell}$$

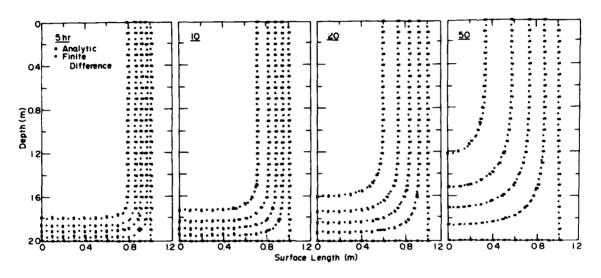
$$\cdot \sum_{m=0}^{\infty} \frac{(-1)^m}{2m+1} \exp \frac{-\alpha (2m+1)^2 \pi t^2}{4b^2} \cos \frac{(2m+1)\pi x}{2b} .$$
 (71)

A short FORTRAN program, INFSUM, was written to compute for this equation the various space and time increments; it is listed with the output in Appendix A. For the comparison, the rectangle is indicated by $\ell = 1$ m, $\ell = 2$ m, $\ell = 40$ °C, $\ell = 20$ °C, and the thermal diffusivity $\ell = 1$ 0°C is set equal to $\ell = 1$ 1.

ADI was run several times to demonstrate the effect of internodal spacing on the accuracy of the solution. Because of the symmetry of the problem, we have to model only one quarter of the problem, assigning a zero flux boundary condition to edges of the grid that fall on lines of symmetry. The region modeled is that portion of the rectangle which lies in the fourth quadrant of the Cartesian graph. The first run was made with an internodal spacing of 0.1 m and a time step of 5 hr. The resulting 24°, 28°, 32°, 36° and 40°C isotherms are plotted for several time steps (Fig. 15a). An inspection of the data printed in ADITMP reveals that the maximum discrepancy between the two solutions is 1.5°C, and it occurs in the first time step (5 hr) in the lower right-hand corner. This location at early times represents the steepest temperature gradient in the problem. By the second time step (10 hr), the discrepancy reduces to a maximum of 0.6°C. The maximum difference between the solutions continues to decrease as the temperature gradient decreases



a. Internodal spacing 0.1 m; time step 5 hr.



b. Internodal spacing 0.05 m; time step 2.5 hr.

Figure 15. Locations of 24°, 28°, 32°, 36° and 40°C isotherms (temperature increasing from left to right) from ADI (finite difference) and analytical solution.

and the location of the maximum temperature difference changes to the locations of the steepest gradients. The input and output for this run of ADI are listed in Appendix A.

The same problem was run with an internodal spacing of 0.05 m and a time step of 2.5 hours. The resulting isotherm plots are shown in Figure 15b. A closer agreement was found between the results of ADI and the analytical solution. A node-by-node comparison shows the maximum difference to be 1°C for the 5-hr distribution, decreasing to 0.4° in the 10 hr distribution. In general, the accuracy of the solution is improved as the internodal spacing decreases; the user must determine the accuracy demanded.

One-dimensional semi-infinite problem

It is also of interest to compare the results of ADI to the one-dimensional problem of a medium initially at a uniform temperature; the surface temperature then undergoes a step change to a different temperature, and the resulting temperature distribution is examined over time. This



Figure 16. One-dimensional semi-infinite problem.

problem corresponds to the theoretical setup of the experimental problem examined in the Comparison of ADI with Experimental Results section. The problem is illustrated in Figure 16.

The solution to this analytical problem is found in many texts on heat transfer; the reader may refer to Holman (1972). The solution is given as

$$\frac{T(x,t)-T_0}{T_i-T_0} = \operatorname{erf} \frac{x}{2\sqrt{\alpha t}}$$
 (72)

where x = the distance below the surface

t = time

 α = the thermal diffusivity, 0.00251 m²/hr in this example.

The grid used for this was 3 by 40 nodes, with an internodal spacing of 0.025 m and a time step of 0.25 hr.

For two depths, 0.05 and 0.20 m, the results of this analytical solution are compared to the results of ADI over 9 hr. The results are given in Table 1.

Table 1. Comparison of ADI results to analytical results for the one-dimensional semi-infinite problem.

x (m)	t (hr)	T (x, t)	ADI
0.05	0.5	26.42	26.66
	1.0	29.12	29.16
	2.0	31.42	31.41
	3.0	32.51	32.51
	4.0	33.18	33.18
	5.0	33.65	33.65
	6.0	34.00	34.00
	7.0	34.27	34.27
	8.0	34.50	34.49
	9.0	34.68	34.68
0.20	0.5	21 11	21.12
	1.0	21.19	21.21
	2.0	21.88	21.90
	3.0	22.83	22.85
	4.0	23.75	23.76
	5.0	24.56	24.56
	6.0	25.26	25.27
	7.0	28.88	25.88
	8.0	26.42	26.42
	9.0	26.89	26.89
			20.07

The results compare to within 0.02°C for times after 2 hr. Earlier times have temperatures that differ by as much as 0.24°C near the top of the grid. Probably, this is because the steepest temperature gradients occur in this problem soon after the step change in temperature occurs at the surface. Recall that steeper temperature gradients require a smaller internodal distance; if we want greater accuracy, the model could be again run using an internodal spacing of less than 0.025 m. The smaller the internodal distance, the more accurate the solution.

Comparison of ADI with experimental results

As an example of the use of ADI, let us compare it to some experimental results. Data are available from a full-scale experiment that was conducted in a 4- by 5-m container of soil that was 1.2 m deep. Of the 4-m width, only the center 1.76 m was included in the experiment; the rest was outside an insulated boundary and was put there for support. The soil used was Lebanon sand of 19.7% moisture content. The experiment was designed to start at an initial uniform temperature of 70°F (21.11°C) in the sand in the box, then the surface temperature was to be raised 100°F (37.78°C), and the change in the temperature distribution over time was to be monitored. An analytical solution is available for this case, and ADI was also run for this case. The results compare very well (see the *One-Dimensional Semi-infinite Problem* section, analytical comparisons).

In the comparison of the experimental data to calculated results, temperatures were used that represent a 20-cm wide band taken vertically through the center of the box. Three strings of thermocouples were placed in this band, as shown in Figure 17.

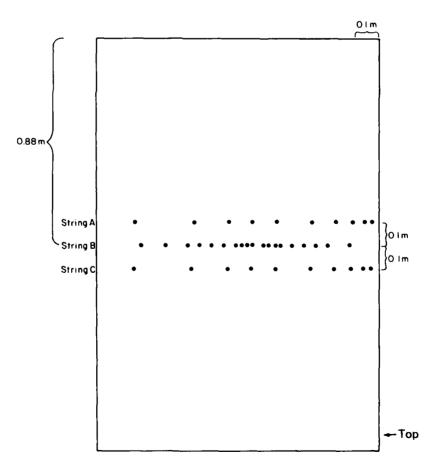


Figure 17. Thermocouple locations in the experimental setup (top of soil container on right).

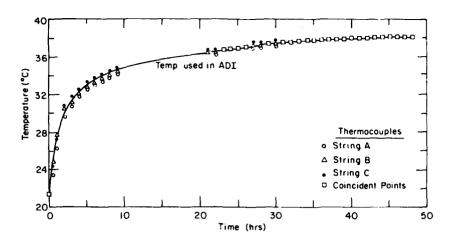


Figure 18. Surface temperatures from experiment used in ADI runs.

In the actual experiment, the initial temperatures ranged from 20.6° to 22.4°C. The initial temperature distribution for ADI was taken as an average of the three thermocouple string temperatures for each depth. It required approximately 27 hr to raise the surface temperature to 37.78°C, after which the surface temperature remained in the range from 37.2° to 38.4°C. Because the data logger malfunctioned, 12 hr of data are missing in a period beginning 9 hr after the start of the experiment; however, the rest of the equipment continued to operate normally. In comparisons made with these data it is assumed that the surface temperature increased linearly over time for this period. See Figure 18 for a graph of the surface temperature vs time for the data and that used in the comparison run with ADI.

Soil tests were conducted at CRREL to determine the density and thermal conductivity of the sand. The density was determined to be 1996.25 kg/m³; this density was assumed constant in the ADI comparison. The conductivity was measured at two temperatures above freezing. At 4.44°C the conductivity was determined to be 1.673 W/m K, and at 26.67°C it was 1.803 W/m K. Each node in the ADI model was assigned a conductivity according to its temperature (in degrees C) for each time step from the following equation:

$$k_{i,i} = 0.00585T + 1.647$$
 (73)

This equation was determined from the two conductivity tests.

The value of the specific heat was taken from measurements done on Lowell sand by Kersten (1949). Again, each node in the ADI model was assigned a temperature-dependent specific heat. The equation fit to Kersten's data is

$$C_{\mathbf{P_{i,i}}} = 0.0039T + 0.34927. \tag{74}$$

The 2 by 48 grid used in ADI had a distance between nodes of 0.025 m. The temperature at the surface was specified, and was taken from the data previously discussed. The time step was 0.25 hr; the surface temperature for each time step was interpolated linearly from the data. The sides and bottom of the grid were assigned the zero heat flux boundary condition. The problem was modeled for 48 hr.

The graphs in Figure 19 compare the results for several times (9 hr, 24 hr and 48 hr) during the run; they show a reasonable agreement but the model tended to underpredict the temperature change.

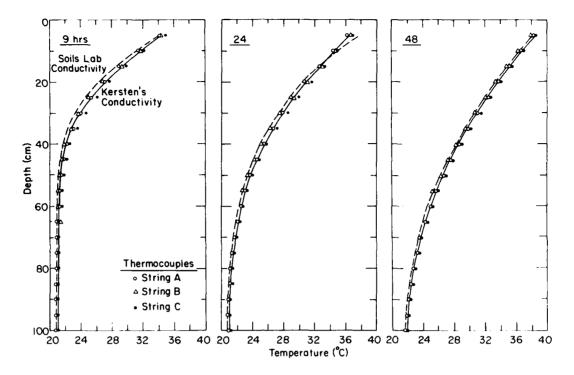


Figure 19. Results of two ADI runs using different conductivities plotted against experimental results.

The model was then run again, using conductivity measurements on Lowell sand by Kersten (1949). The temperature-dependent equation for conductivity used for this run was

$$k_{i,i} = 0.00585T + 1.975$$
 (75)

These results may again be found in Figure 19. This time excellent agreement was found between the actual data and the ADI calculations.

Each of the two runs of ADI required 2 minutes and 13 seconds of computer time on the PRIME computer at CRREL.

VERIFICATION OF ADIPC

Comparison of ADIPC with analytical results-the Neumann solution

ADIPC was first compared to the well-known one-dimensional analytical Neumann solution. A semi-infinite region is initially at a uniform temperature T_0 which is above the fusion temperature $T_{\rm f}$. Suddenly, the surface temperature is changed to $T_{\rm s}$, a temperature below the fusion temperature. The subsequent movement of the phase change front may be calculated.

Allow ρ , C_P , k and H_L to represent the density, specific heat, thermal conductivity and latent heat of fusion, respectively. Thermal diffusivity α is defined as $k/\rho C_P$. Temperature T is a function of position and time. Subscripts 1 and 2 refer to the property or variable in the frozen and liquid phase, respectively. The boundary conditions follow:

$$T_1(0,t) = T_s$$
 (76)

$$T_2\left(x,0\right) = T_0 \tag{77}$$

$$T_1(X, t) = T_2(X, t) = T_f$$
 (78)

$$k_1 \frac{\partial T_1}{\partial X} - k_2 \frac{\partial T_2}{\partial X} = \rho H_L \frac{\partial X}{\partial t}. \tag{79}$$

The temperatures in the liquid and solid regions must satisfy the following equations

$$\frac{\partial^2 T_1}{\partial x^2} - \frac{1}{k_1} \frac{\partial T_1}{\partial x} = 0 \tag{80}$$

$$\frac{\partial^2 T_2}{\partial x^2} - \frac{1}{k_2} \frac{\partial T_2}{\partial x} = 0. \tag{81}$$

The solution has been given many times in the literature (e.g. Carslaw and Jaeger 1959). The solution for the location of the phase change front follows:

$$X = 2\lambda \sqrt{\alpha_1 t} \tag{82}$$

where λ must be determined from the relation

$$\frac{e^{-\lambda^2}}{\operatorname{erf} \lambda} \frac{k_2 \sqrt{\alpha_1} (T_0 - T_f) e^{-\alpha_1/\alpha_2 \lambda^2}}{k_1 \sqrt{\alpha_2} (T_f - T_s) \operatorname{erfc} \lambda \sqrt{\alpha_1/\alpha_2}} = \frac{\lambda H_L \sqrt{\pi}}{C_{P_1} (T_f - T_s)}.$$
 (83)

For comparison the following values were used:

$$k_1 = 2.21 \frac{W}{m K}$$
 $k_2 = 0.580 \frac{W}{m K}$
 $\rho_1 = 917 \frac{kg}{m^3}$ $\rho_2 = 998.2 \frac{kg}{m^3}$
 $C_{P_1} = 0.5815 \frac{W hr}{kg K}$ $C_{P_2} = 1.16 \frac{W hr}{kg K}$
 $H_L = 93.00 \frac{W hr}{kg}$
 $T_f = 0^{\circ}C$.
 $T_s = -4.67^{\circ}C$

The problem was modeled for 3 hr using an internodal spacing of 0.5 cm and a time step of 0.0025 hr. The depth of the 0°C isotherm is plotted with the analytical solution in Figure 20. As mentioned earlier, the location of the phase change front is found by interpolating between the nodal temperatures to find the 0°C isotherm. The front progresses in a step-like pattern. This occurs when the location of phase change moves from one node to the next and is inherent in the apparent heat capacity method in a discretized space. Nevertheless, the results of ADIPC show good agreement with the analytical solution. A more accurate computed solution could be obtained by using a smaller internodal spacing and smaller time step. A copy of ADIPC and output for this solution is included in Appendix B.

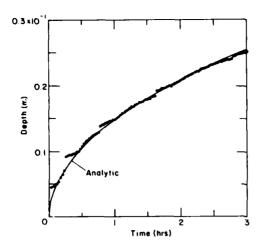


Figure 20. Comparison of ADIPC with Neumann solution.

Comparison of ADIPC with analytical results-two-dimensional phase change verification

Next compare the results of ADIPC with an analytical solution involving phase change around a pipe. This problem is two-dimensional in Cartesian coordinates and thus is a two-dimensional verification of ADIPC. But it is a one-dimensional problem in cylindrical coordinates, facilitating an analytical solution. An exact analytical solution is available for the case of freezing in a region, initially at a uniform temperature, that is suddenly subject to the effects of a continuous line source that extracts heat at a rate of Q per unit time. The exact solution of this heat conduction problem is given in Carslaw and Jaeger (1959). The results are as follows.

The location of the freezing front at any time is given by

$$R = 2\lambda \sqrt{\alpha_1 t} \tag{84}$$

where R is the radius of the front (r = 0) is the location of the line source), α_1 is the thermal diffusivity of the frozen zone, t is time and λ is obtained from the following relation:

$$\frac{Q}{4\pi} \exp(-\lambda^2) + \frac{k_2 (T_i - T_f)}{E_i (-\lambda \alpha_1 / \alpha_2)} \exp(-\lambda^2 \alpha_1 / \alpha_2) = \lambda^2 x_1 L\rho$$
 (85)

where k =conductivity

 α = diffusivity

 T_i = initial temperature

 $T_{\rm f}$ = fusion temperature

 $L\rho$ = latent heat per unit volume.

 E_i is the exponential intergral function

$$E_i(x) = \int_{-\infty}^{x} \frac{e^{\nu} dV}{V}.$$

and the subscripts 1 and 2 represent the frozen and unfrozen zones, respectively.

The temperatures in the frozen and unfrozen zones are given by

$$T_1 = T_f + \frac{Q}{4\pi k_1} \left[E_i \left(-\frac{r^2}{4\alpha_1 t} \right) - E_i \left(-\lambda^2 \right) \right] \qquad 0 < r < R$$
 (86)

$$T_2 = T_i - \frac{(T_i - T_f)}{E_i (-\lambda^2 \alpha_1/\alpha_2)} \quad E_i \left(\frac{-r^2}{4\alpha_2 t}\right) \qquad r > R.$$
 (87)

In order to get a solution for phase change around a pipe from this, the pipe is assigned a constant radius r_p ; the temperature at r_p varies with time according to eq 86. A shifted time is used so that at time t = 0 in the computer run, the location of the phase change front is at $r = r_p$. For this time, the initial temperature distribution is figured from eq 87.

For the comparison, the following values were used:

$$k_1 = 0.0072 \frac{\text{cal}}{\text{cm s °C}}$$
 $k_2 = 0.0042 \frac{\text{cal}}{\text{cm s °C}}$
 $\alpha_1 = 0.014165 \frac{\text{cm}^2}{\text{s}}$
 $\alpha_2 = 0.005556 \frac{\text{cm}^2}{\text{s}}$
 $L\rho = 33.012 \frac{\text{cal}}{\text{cm}^2 °C}$
 $T_i = 4 °C$
 $T_f = 0 °C$.

Densities for all regions were assumed equal in the solution.

Equation 84 was solved by substituting a polynomial approximation for the exponential integrals (Abramowitz and Stegun 1970) and then solving the equation for λ by an iterative scheme on the computer. For this case it was found that $\lambda = 0.08246$.

ADIPC used an internodal distance of 1 cm and a time step of 60 s on a 30 by 30 grid. Because of the symmetry of the problem, with the Cartesian origin at the center of the pipe, only one-quarter of the situation was modeled. The top and right-hand edges of the grid were assigned zero flux boundaries, and the left-hand and bottom edges of the grid were assigned the semi-infinite condition.

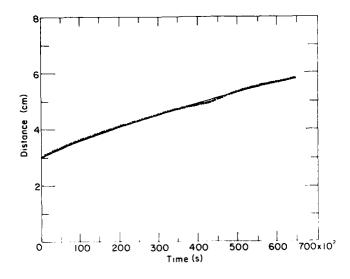


Figure 21. Comparison of ADIPC with radial analytic freezing solution.

The location of the freezing front over time was found in ADIPC by interpolating to find the zero degree isotherm; this is plotted against the analytical solution in Figure 21. Excellent agreement is found.

USER INSTRUCTION FOR ADI

The first step in using ADI to solve a problem is to define the boundaries of the problem and to identify the different materials in the problem. It is important to know the dimensions of objects to be modeled as accurately as possible. Next, identify the boundary conditions in the problem. Now draw the grid, determining a nodal spacing that will accurately represent the materials in the problem. Indicate the boundary conditions on the grid, then set these conditions in the computer program by assigning the appropriate values of RAY (I, J, 2) to each node I, J in the grid; the values are listed in the comment statements at the beginning of subroutine ADDATA.

Material properties such as thermal conductivity, density, etc., must be known in consistent units. There are no dimensional constants inherent in the program, so any system may be used. For the problems presented in this report, the following units were used: thermal conductivity (W/m K), density (kg/m³), specific heat (W hr/kg K), temperature (°C), time (hr) and distance (m). The units of the convection coefficient would be W/m²K. Values for the conductivity, specific heat, density and convection coefficient must be specified by the user as indicated in the main program (near the start of the 2001 and 2002 loops). These values may be programmed to change with temperature or time and could conceivably be different for each node in the grid. At the time of this report's publication, the conductivity of a semi-infinite node in the program is assumed equal to that of the adjacent regular interior node.

Now edit subroutine ADDATA following the directions in the comment statements in the subroutine to initialize the variables and arrays. Note that the variables and arrays are defined at the beginning of ADDATA.

When editing subroutine ADDATA, the user will encounter a variable named ITRT, which is defined as the number of time steps before the results are printed. Setting this value as five, for example, will set a counter in the main program so that at every fifth complete time step the main program will print the temperature distribution (RAY [I, J, 1]) into file ADITMP. Similarly, variable ITPC will set a counter to call subroutine ISOTHM to locate the user-specified isotherms in the current temperature distribution. The coordinates of these isotherms are printed into file POINTS.

When ADI is run, it will interactively ask whether or not the user wishes to run subroutine ADDATA. The first time the program is run, ADDATA must be run. The variables will be put into a formatted data file, ADIDAT. If, after the program is run, a change is made in the main program but not to ADDATA, the subroutine need not be called again; the main program will read the data from ADIDAT.

Once the program has been run, the user should examine output file ADIOUT closely to be sure that the initial conditions and boundary conditions are those intended, i.e., that no mistake was made when editing ADIDAT or changing values in the main program.

USER INSTRUCTION FOR ADIPC

The use of ADIPC is the same as ADI with the exception of the specification of several phase change variables. When editing ADDATA, the user must specify the value of the latent heat of fusion, the conductivity and density of the phase change region, and the temperature range over which phase change will occur. If the units indicated in the instructions for ADI are used, the

latent heat will have the units W hr/kg, as does the specific heat. The value of the apparent specific heat (discussed in the *Phase Change* section), known as CPPC in the program, should be calculated by the user as follows:

$$CPPC = \frac{1}{2} \left(C_{\mathbf{p}_{\mathbf{B}}} + C_{\mathbf{p}_{\mathbf{A}}} \right) + \frac{H_{\mathbf{L}}}{\Delta T}$$

where

 $C_{P_{R}}$ = specific heat of the frozen material

 $C_{\mathbf{P}_{\mathbf{A}}}$ = specific heat of the unfrozen material

 $H_{\rm L}$ = latent heat of fusion for the material

 ΔT = temperature range around the fusion temperature where phase change occurs from

$$T_{\rm f} - \frac{\Delta T}{2}$$
 to $T_{\rm f} + \frac{\Delta T}{2}$.

 ΔT in the verifications was 1.0°C.

As in ADI, the values for density, specific heat and conductivity must be specified as indicated in the comments near the start of the main program, ADIPC.

CONCLUSIONS

Two two-dimensional finite difference computer programs have been developed to model time-varying heat conduction. Results of test runs of the programs show excellent agreement with analytical and experimental results. The programs are easily set up to model new problems, and have the capability to solve a wide variety of heat conduction problems.

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APPENDIX A: PROGRAM INFSUM AND SAMPLE INPUT AND OUTPUT FOR PROGRAM ADI

ৰ প্ৰথম কৰিবলৈ কৰিবলৈ কৰিবলৈ কৈ লোক কৰিবলৈ জানুক্ত কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিব সংস্থানিক কৰিবলৈ কৰিবলৈ কৰিবলৈ কৈ কৰিবলৈ কৰিবলৈ জানুক্ত কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবলৈ কৰিবল

INFSUM, a short program to do infinite sum calculation for rectangle problem

```
INFSUM
THIS CALCULATES THE IMPINITE SUM PRODUCT FOR THE RECTANGLE.
IMPLICIT DOUBLE PRECISION(A-H.O-Z)
IMPLICIT INICGER+4(1-N)
DIMENSION TEMP(50.05)
CALL CONTRL(2.*INFOUT*.5)
        Mab:
A=.00251
EL=1
d=0
T0=40.
         TI=20.0
PI=3.141592654
PIC=9.869654464
       FIGURE SUM2
N=1
SUM2=0.
DO 230 INN=1.50
IN=INN-1
C3A=C3*(2.*IN+1.)*(2.*IN+1.)
C4A=C4*(2.*IN+1.)
C3A=DEXP(C3A)
SCND=(N/(2.*IN+1.))*C3/A*DCOS(C4A)
SUM2=SUM2+SCNO
N=-1*N
IF(C3AA.LE...)CO1) GU TO 231
      SUM2=SUM2+SCNO
V=-1*N
IF(C3AA.LE..3C01) G. TO 231
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
GONTINUE
GONTINUE
FORMAT(/.1X..TEMPERATURES AFTER..F6.2.. HOURS.)
WRITE(5.212) ((TEMP(I.J).J=1.IX).I=1.IY)
CONTINUE
CONTINUE
CALL CONTRL(4..INFOUT..5)
CALL EXIT
       ĔND
```

Sample output from program INFSUM

```
TEMPERATURES 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
20.000 20.000
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TEMPERATURES WILL BE PRINTED EVERY 1 TIME STEPS.

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LEFT BOUNDARY
CONSTANT HEAT FLUX
BOTTOM LEFT CORNER CONSTANT TEMPERATURE
BOTTOM EGINDARY
CONSTANT TEMPERATURE
BOTTOM RIGHT CORNERCONSTANT TEMPERATURE
RIGHT BOUNDARY
CONSTANT TEMPERATURE RIGHT CORNER CONSTANT TEMPERATURE

NUMBER OF TIME STEPS=

Sample output from program ADI for rectangle problem (file ADITMP)

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32.30
34.28
37.03
```

Sample of isotherm locations output from ADI for rectangle problem (file POINTS)

```
THE FOLLOWING 144 POINTS REPRESENT TIME = 5.00 HOURS:

40.00 DEGREE ISOTHERM:
0.000 0.200 2.0000
0.200 2.0000
0.300 2.0000
0.400 2.0000
0.500 2.0000
0.600 2.0000
0.600 2.0000
0.600 2.0000
0.600 2.0000
0.900 2.0000
1.000 0.2000
1.000 0.2000
1.000 0.5000
1.000 0.5000
1.000 0.5000
1.000 0.5000
1.000 0.5000
1.000 0.5000
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APPENDIX B: PROGRAM ADIPC AND SAMPLE INPUT AND OUTPUT

Program ADIPC and its subroutines

```
23345677899011233451
                                                                                    ADIPC
                   THIS FORTRAN PROGRAM SOLVES FOR TWO DIMENSIONAL TEMPERATURE DISTRIBUTION RESULTING FROM CONDUCTION HEAT TRANSFER. BOUNDARY CONDITIONS MAY INCLUDE SPECIFIED TEMPERATURES, CONVECTIVE SURFACES, AND SEMI-INFINITE BOUNDARIES. ADIPC USES AN IMPLICIT ALTERNATING DIRECTION FINITE DIFFERENCE NUMERICAL TECHNIQUE ON A GRID WITH SQUARE ELEMENTS. THE MATERIAL PROPERTIES (TK.CP.RO.H) MAY BE DIFFERENT FOR EACH NODE OF THE GRID. AND MAY CHANGE WITH TIME. ADIPC USES THE APPARENT HEAT CAPACITY METHOD FOR PHASE CHANGE. THIS PROGRAM WAS WRITTEN BY MARY REMLEY ALBERT AT CRREL. 1981.
DATA FOR ADJPC IS GATHERED BY ADDATA, WHICH PUTS IT INTO FILE ADPDATA SEE ADDATA FOR AN EXPLANATION OF VARIABLES, AND TO SET UP THE PROGRAM FOR YOUR PROBLEM.

ADJUSTMENTS FOR CHANGING MATERIAL PROPERTIES WITH TIME (DURING THE RUN) SHOULD BY TYPED INTO THE MAIN PROGRAM AS INDICATED NEAR LINE 170.
                 ----
                    38)
39)
40)
41) C
42)
43)
44) C
45)
46)
46)
48)
49)
50)
             1
51)
52)
53)
54)
55)
56) 4
57) C
             10
61)
             20
64)
              15
 66)
67)
              30
68)
70)
71)
72)
73)
75)
              35
               34
```

```
READ(5,35) (TMPR(I),I=1,Y)

CONTINUE

WRITE INITIAL DATA INTO OUTPUT FILE ADPOUT

WRITE (6,21)

FORMAT(/,1X,*DATA FOR THIS RUN OF ADIPC:*,/)

WRITE(6,22)

FORMAT(/,1X,* DS DELT DI TDEL*)

WRITE(6,20)

WRITE(6,23)

FORMAT(/,5X,*A*,4X,*X*,4X,*Y*,NISO ITRT IMAX ITPC *)

WRITE(6,20)

WRITE(6,20)

WRITE(6,20)

WRITE(6,20)

FORMAT(/,1X,* ROPC(K) CPPC(K) HL(K) TPC(K)*)

WRITE(6,20)

FORMAT(/,1X,*TISO(B),B=1,NISO:*)

WRITE(6,20)

FORMAT(/,1X,*TISO(B),B=1,NISO:*)

WRITE(6,20)

FORMAT(/,1X,*TISO(B),B=1,NISO:*)

WRITE(6,25)

FORMAT(/,3X,*H(K)*)

WRITE(6,30) (H(L),L=1,A)

WRITE(6,30) (H(L),L=1,A)

WRITE(6,30) (KRNR(J),J=1,4)

FORMAT(/,1X,*KRNR(1)=*,12,* KRNR(2)=*,12,* KRNR(3)=*,12,*

WRITE(6,26)

FORMAT(/,3X,*FLXT(J)*,3X,*FLXB(J)*,3X,*TMPT(J)*,3X,*

FORMAT(/,3X,*FLXT(J)*,3X,*FLXB(J)*,3X,*TMPD(J)*,3X,*

FORMAT(/,3X,*FLXT(J)*,3X,*FLXR(I)*,3X,*TMPD(I)*,3X,*

FORMAT(/,3X,*FLXL(I)*,FLXR(I)*,3X,*TMPD(I)*,3X,*

FORMAT(/,3X,*FLXL(I)*,5XR(I)*,3X,*TMPD(I)*,3X,*

WRITE(6,27) ((FLXL(I),FLXR(I)*,TMPD(I)*,TMPR(I)*,1=1,Y)

RXX=X/17

JJX=RX

JJR=X-(1,7*JJX)
        81)
82)
83) c<sup>12</sup>
84)
85) 21
86)
87) 22
           88)
           90)
91)
92)
93)
95)
                                                             23
                                                              29
          96)
97)
28)
                                                             24
 95)
                                                              25
101)
162)
103)
                                                              33
104)
105)
106)
107)
108)
110)
111)
112)
                                                             26
                                                             27
                                                              28
                                                                                                            ## THERE(1)*)

WRITE(6,27) ((FLXL(I),FLXR(I),TMPL(I),TMPR(I)),I=1,Y)

RX=X/17

JUX=RX

JUX=X-(17*JJX)

DO 63 K=1-3

GO TO (64,65,66)*K

WRITE(6,31)

GO TO 67

WRITE(6,45)

GO TO 67

WRITE(6,44)

DO 68 JJ=1*JJX

IF(JJX,EQ.0) GO TO 68

J2=17*JJ

J1=J2-16

WRITE(6,500) ((RAY(I,J,K),J=J1,J2),I=1,Y)

IF(JJX,LI,JJX) WRITE(6,503)

CONTINUE

J2=X

J1=X-JJR+1

IF(JJX,EQ.0) GO TO 69

IF(JJX,NE-0) WRITE(6,503)

DO 69 I=1,Y

WRITE(6,500) (RAY(I,J,K),J=J1,J2)

CONTINUE

CONTINUE

CONTINUE

CONTINUE

FORMAT(/,1X,*RAY(I,J,2) NODAL MATERIAL TYPE*)

FORMAT(/,1X,*RAY(I,J,2) NODAL LOCATION TYPE*)
114)
115)
116)
117)
118)
119)
120)
122)
1223)
1224)
1226)
1227)
1228)
1230)
131)
132)
                                                             64
                                                              65
                                                             68
CCCCC
13378)
113390)
114423)
11444678)
11444678)
115567)
115589)
115589)
                                                             69
63
44
45
31
                                                             11
                                                              37
                                                                                               161) C
162) C
163)
164)
165)
166)
                                                                                                                   INITIALIZE SPECIFIC HEAT AND CONDUCTIVITY
IF TEMPERATURE DEPENDENT, INSERT EQUATIONS.
DO 46 J=1.X
DO 47 I=1.Y
TEMP(I.J)=RAY(I.J.1)
K=RAY(I.J.3)
IF(RAY(I.J.1).LT.TPC(K)) GO TO 39
UNFROZEN
  168) C
169)
170)
171)
172)
                                                                                                                   UNFRUZEN
CP(I+J)=1.160D0
RO(I+J)=998.20D0
GO TO 47
CONTINUE
FROZEN
    173) c<sup>39</sup>
```

```
CP(I.J)=.5815D0
RO(I.J)=917.000
CONTINUE
       1/9;
180;
181;
182;
183;
                                                                                                                                                                                                                                               COMES HERE AFTER EVERY COMPLETE TIME STEP NOTE THAT Q=0 FOR 1ST ITERATION
  181)
184)
184)
184)
186)
1887)
1887)
1891)
1991)
1992)
1993)
1995)
                                                                                                         IF SPECIFIED TEMPERATURES CHANGE WITH TIME. SPECIFY THE CHANGES HERE. (NUMBERS 57 AND 58 ARE FREE FOR LOOPS IF NEEDED.)
                                                                                          ADJUST THERMAL PROPERTIES WITH TEMPERATURE.

IF APPROPRIATE, USE RAY(I,J.3) TO INDICATE MATERIAL TYPE.

(BUT DO NOT CHANGE THE NEXT 9 LINES.)

DO 50 J=1.4X

NO 51 I=1.4Y

     195)
196)
197)
198)
199) C
        200)
        201)
202)
203)
203) 60
204) C
205)
206)
207)
2081
                                                                                                                                                                                                                                   FROZEN
                                                                                                                              FROZEN

ISTAT(1.J)=1

TK(1.J)=2.210D0

RO(1.J)=917.0D0

CP(1.J)=.58150D0

GO TO 51

CONTINUE
  $\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{2}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2
     41
                                                                          161
                                                                          42
                                                                                                                             CONTINUE
FIGURE RESULTANT CONDUCTIVITIES BETWEEN NODES
KONT=KONT+1
INDEX=RAY(I = J = 3)
If (I = EQ = 1) GO TO 71
IK1 = (2 = DDO*TK(I + J)*TK((I - 1) + J))/(TK(I + J) + TK((I - 1) + J))
CONTINUE
IF (J = EQ = 1) GO TO 72
IK2 = (2 = CDO*TK(I + J) + TK(I + (J - 1)))/(TK(I + J) + TK(I + (J - 1)))
CONTINUE
IF (I = EQ = Y) GO TO 73
IK3 = (2 = DDO*TK(I + J) + TK((I + 1) + J))/(TK(I + J) + TK((I + 1) + J))
CONTINUE
IF (J = EQ = X) GO TO 74
IK4 = (2 = DDO*TK(I + J) + TK(I + (J + 1)))/(TK(I + J) + TK(I + (J + 1)))
CONTINUE
IF (I = EQ = X) GO TO 75
IF (I = EQ = X) GO TO 75
NODE PRESENTLY PHASE CHANGE
                                                                          1601
                                                                        71
                                                                        72
                                                                          73
                                                                        74
        266)
267) C
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282)
283)
                                                                                                                                        R3=2.DC*DS*DS*ROPC(INDEX)*CPPC(INDEX)/DELT
R4=R3
GO TO 7B
CONTINUE
NODE NOT PHASE CHANGE PRESENTLY
R3=2.DC*DS*DS*ROO(I.J)*CPO(I.J)/DELT
R4=2.DC*DS*DS*RO(I.J)*CP(I.J)/DELT
CONTINUE
ROUTE CORNERS
                                                                                                                             CONTINUE

ROUTE CORNERS

IF(RAY(I,J,2),GE.5) GO TO 38

IJ2=RAY(I,J,2)

GO TO (101.102.103.104).IJ2

ROUTE THE REST OF THE NODES

CONTINUE

IJ2=RAY(I,J,2)-4

GO TO (105.106.107.108.109.110.111.112.113.114.115.

$ 116.117.118.119.120.121.122.123.124.125.126.127.

$ 128.129.130.131.132).IJ2
                                                              284)
285)
285)

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320)
319) 1104 E

320) E

321) E

323) B

323) B

323) C

326) 1101

327) C

328) I

329) I

3331) C

3342) C

3442) C

3443) C

3443) C

3443) C

3445) C

3446) C

3446) C

3447) C

3457) C

347) C

347
   345)
346)
347)
348)
350) 109
351) C
352) C
353)
353)
355)
     356)
357)
358)
359)
360)
361)
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PRODUCTION OF SECURITY BUSINESS OF MANAGEMENT

```
362)
363)
EB(KONT)=-TK3-.5D0+R4
364)
C(KONT)=TK3
365)
D(KONT)=-SD0+TK2+RAY(I.(J-1).1)-.5D0+TK4+RAY(I.(J+1).1)+
366)
L(SD0+TK2+.5D0+TK4-.5D0+R3+H(INDEX)+DS)+RAY(I.J-1)-H(INDEX)+DS+
367)
GD TO 165
369)
GD TO 165
370)
C
IFLAG(7)=12
ARIGHT SIDE CONVECTIVE
371)
372)
IF(JFLG.NE.1) GO TO 407
374)
EA(KONT)=TV0
```

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GO TO 165
 456)
457)
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459)
    496)
 498)
499) C
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531)
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| Top LEFT | Top |
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MANAGER VENEZIA NASSESSA LEGISLADO

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120 CONTINUE

TOP CONST FLUX NODE ADJ TO RIGHT SEMI-INF

IF (JFLG.NE.1) GO TO 1200

EA(KONT)=-50D0*TK2

EB(KONT)=--50D0*(TK2+TK4/DI+R4)

C(KONT)=TK4/(2.0D0*DI)

D(KONT)=-TK3*RAY((I+1).J.1)-DS*FLXT(J)+(TK3-.50D0*R3)

*RAY(I.J.1)

GO TO 165

1200 CONTINUE

EA(KONT)=0.0D0

EB(KONT)=-TK3-.50D0*R4

C(KONT)=TK3

D(KONT)=TK3

D(KONT)=-1.*(TK4/(2.0D0*DI))*RAY(I.(J+1).1)-.5*TK2
  721)
722)
723)
724)
725)
726)
728)
728)
731)
733)
              C
                              D(KONT)=-1.*(TK4/(2.0D0*DI))*RAY(I,(J+1),1)-.5*TK2
```

ELECTRICAL PROPERTY ASSESSED ASSESSED

```
**RAY(I, (J-1), 1)*(TK4/(2.000*DI)*.5000*TK2-.5000*R3)

**RAY(I, J-1)-0S*FLXT(J)*

CONTINUE

IOP CONST FLUX ADJAC TO LEFT SEMI-INF

IF(J-EG.ME.1) GO TO 1220

EA(KONT)=-K2/(2.000*DI)*

EE(KONT)=-K300*TK2/DI-TK4+R4)

CKONT)=-S000*TK2/DI-TK4+R4)

CKONT)=-S000*TK2/DI-TK4+R4)

CKONT)=-S000*TK3/CI-J-I)-DS*FLXT(J)*

**(TK3--5000*R3)*RAY(I, J, I)-DS*FLXT(J)*

**(TK3--5000*R4)*

**(TK3--5000*TK4)*

**(
 734)
735)
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738)
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753) 114
754) 118
755) 124
756) 125
757) 126
758) 760) C
761) C
763) C
764) C
765) 202
766) C
767) 767)
7769) 7770)
7771)
7772)
7773)
7774)
                                                     c<sup>202</sup>
 203 CONTINUE

CONVECTIVE ON BOTH SIDES

IFLAG(2)=2

IF(JFLG.NE.1) GO TO 2030

EA(KONT)=0.D0

EB(KONT)=-.5D0*TK4-.25D0*R4

C(KONT)=-.5D0*TK3*RAY((I+1),J.1)+(.5D0*TK3+H(INDEX)*DS-.25D0*R3)*

8 RAY(I,J,1)-.5D0*H(INDEX)*TMPT(J)-.5D0*H(INDEX)*TMPL(I)

GO TO 1.65

2030 CONTINUE

EA(KONT)=-.5D0*TK3-.25D0*R4

C(KONT)=-.5D0*TK3-.25D0*R4

C(KONT)=-.5D0*TK4*RAY(I,(J+1),1)+(.5D0*TK4+H(INDEX)*DS-.25D0*R3)*

8 RAY(I,J.1)-.5D0*TK4*RAY(I,(J+1),1)+(.5D0*TK4+H(INDEX)*DS-.25D0*R3)*

8 RAY(I,J.1)-.5D0*H(INDEX)*(TMPT(J)+TMPL(I))

GO TO 1.65
   789)
7791)
7792)
7793)
7795)
7796)
7799)
                                                                                                                          & RAY(I.J.1)
GO TO 165
CONTINUE
                                                                     GO TO 165
CONTINUE
SEMI-INF ON BOTH SIDES

IFLAG(2)=3
WRITE(1.9997) I.J
GO TO 9999

206
CONTINUE

VERT CONST FLUX--HORIZ CONVECT

IFLAG(2)=4
IF(JFLG.NE.1) GO TO 2060
EA(KONT)=-50D0*TK4-.250D0*R4
C(KONT)=-50D0*TK4*
D(KONT)=-50D0*TK3*RAY((I+1).J.1)-.50D0*FLXL(I)*DS
*-.50D0*H(INDEX)*DS*TMPT(J)+(.50D0*TK3+.50D0*H(INDEX)*DS
for TO 165
CONTINUE
EA(KONT)=-50D0*TK3-.50D0*H(INDEX)*DS
C(KONT)=-50D0*TK3-.50D0*H(INDEX)*DS
C(KONT)=-50D0*TK3-.50D0*H(INDEX)*DS
C(KONT)=-50D0*TK3-.50D0*H(INDEX)*DS
C(KONT)=-50D0*TK3+.50D0*FLXL(I)*DS
*-.50D0*H(INDEX)*DS*TMPT(J)+(.50D0*TK4-.250D0*R3)

*RAY(I.J.1)
GO TO 165
CONTINUE

TFLAG(2)=7
CONVECT--HORIZ CONST FLUX
                                                                        204
   CONTINUE
VERT CONVECT--HORIZ CONST FLUX
IFLAG(2)=7
IF(JFLG.NE.1) GO TO 2050
EA(KONT)=0.0D0
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828)
829)
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832)
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888) 102
890) C
891)
892) 212
893) C
894)
895)
                                                       GO TO (211-212-213-214-215-216-217-218-219-220) ** KRNR(2)

CONST FLUX ON BOTH SIDES

IFLAG(4)=1

IF(JFLG-NE-1) GO TO 2120

EAKONT)=0-5D0*TK4--25D0*R4

C(KONT)=-5D0*TK4--25D0*R4

D(KONT)=-5D0*TK4-RAY((I-1)-J-1)+(-5D0*TK1--25D0*R3)*RAY(I-J-1)

$ --5*D5*(FLX8(J)+FLXL(I))

GO TO 165

2120 CONTINUE

EAKONT)=-5D0*TK1--25D0*R4

C(KONT)=-5D0*TK1--25D0*R4

C(KONT)=-5D0*TK4-RAY(I-(J+1)-1)+(-5D0*TK4--25D0*R3)*RAY(I-J-1)

$ --5*D5*(FLX8(J)+FLXL(I))

GO TO 165

213 CONTINUE

CONVECT ON BOTH SIDES

IFLAG(4)=2

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1001) C
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1005)
                     218 CONTINUE

VERT=SEMI-INF+HORIZ=CONST FLUX

IFLAG(4)=9

IF(JFLG.NE.1) GO TO 2180

EA(KONT)=0.0D0

EB(KONT)=-(TK4+TK(I,J))/(2.0D0+D1)-.50D0+D12+R4

C(KONT)=TK4/(2.0D0+D1)

D(KONT)=-FLXB(J)+D12+DS-(TK(I,J)/(2.0D0+D1))+TMPL(I)-TK1+D12

**RAY((I-1)*,J*1)+(TK1+D12-.50D0+D12+R3)*RAY(I,J*1)

GO TO 165

2180 CONTINUE

EA(KONT)=TK1+D12

EB(KONT)=TK1+D12

EB(KONT)=TK1+D12-.50D0+D12+R4

C(KONT)=0.0D0

D(KONT)=-(TK4/(2.0D0+D1))*RAY(I,(J+1)*1)-(TK(I,J)/(2.0D0+D1))*
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1012)
1013)
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ANGRESON YPERZHON INGRESSES (INVERSE)

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1025)
                                        # TMPL(I)-FLXB(J)*DI2*DS+((TK4+TK(I,J))/(2.0D0*DI)-.50D0*DI2*R3)
# RAY(I,J,1)
GC TO 165
CONTINUE
VERT=CONVECT.*HORIZ=SEMI-INF
IFLAG(4)=6
IF(JFLG.NE.1) GO TO 2190
EA(KONT)=0.0D0
EB(KONT)=TK4*DI2-H(INDEX)*DI2-.50D0*DI2*R4
C(KONT)=TK4*DI2
D(KONT)=TH(INDEX)*DI2*TMPL(I)-(TK1/(2.0D0*DI3)*RAY((I-1)...In))
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1114)
1116)
                                                                                & TMPR(I)-TK(I,J)+DINF+TMPB(J)+(TK1+DINF+
& 2.00+TK(I,J)+DINF-(DI2+DI2)+R3)+RAY(I,J,1)
GO TO 165
CONTINUE
                                                    2240
                                                                                CONTINUE
EA(KONT)=TK1*DINF
EB(KONT)=-TK1*DINF-(DI2*DI2)*R4
C(KONT)=C.DO
D(KONT)=-TK(1*J)*DINF*TMPR(I)-TK2*DINF*
& RAY(I*(J-1)*1)*(2*D0*TK(I*J)*DINF*TK2*DINF-
& DI2*D12*R3)*RAY(I*J*1)-(TK(I*J)*DINF)*TMPB(J)
GO TO 155
CONTINUE
VERT CONST FLUX-- HORIZ CONVECT
 1117)
1118)
1119)
  1120) C
                                                                                                                          VERT CONST FLUX-- HORIZ CONVECT
                                                                                      VERT CONST FLUX-- HORIZ CONVECT

IFLAG(6)=4

IF(JFLG.NE.1) GO TO 2263

EA(KONT)=-50D0+TK2

EB(KONT)=-50D0*TK2-.250D0*R4

C(KONT)=0.0D0

D(KONT)=-50D0*TK1*RAY((I-1).J.1)-.50D0*FLXR(I)*DS-.50D0*H(INDEX)

B + OS + TMPB(J) + (.50D0*TK1+.50D0*H(INDEX)*DS-.250D0*R3)*RAY(I.J.1)

GO TO 165

CONTINUE
 1121)
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1153)
                                           CONTINUE

CONTINUE

CONTINUE

VERT CONST FLUX, HORIZ SEMI-INF

IFLAG(6)=5.1) 60 TO 2270

E (400NT)=-12.012

E (400NT)=-12.012

E (400NT)=-178.012-.25D0+D12*R4

C (400NT)=-178.112.00*D1)*RAY((1-1)*J.)-(TK(1,J)/(2.D0*D1))*

& TMPB(J)-FLXR(1)*D12*D5*(TR1/(2.D0*D1)*TK(1,J)/(2.D0*D1)-

& 25D0*D12*R3)*RAY(1,J.1)

GO TO 165

2270 CONTINUE

E (4KONT)=-TK1/(2.D0*D1)

E (4KONT)=-TK1/(2.D0*D1)-TK(1,J)/(2.D0*D1)-.25D0*D12*R4

C (KONT)=0.D0

D (KONT)=-TK2*D12*RAY(1,(J-1)*1)-(TK(1,J)/(2.D0*D1))*TMPB(J)-

& FLXR(1)*D12*D5*(TK2*D12*-25D0*D12*R3)*RAY(1,J.1)

C (KONT)=0.D0

D (KONT)=-TK2*D12*RAY(1,(J-1)*1)-(TK(1,J)/(2.D0*D1))*TMPB(J)-

E (4KONT)=-TK2*D12*RAY(1,J.1)/(2.D0*D1)-.50D0*D12*R4

C (50 TO 165

C (50 TO
1154) 227
1155) C
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122220)
122223)
                   1208)
                                                                                 GO TO 9999

236 CONTINUE

VERT CONST FLUX--HORIZ CONVECT

IFLAG(8)=4

IF(JFLG.NE.1) GO TO 2360

EA(KONT)=-.50D0*TK2

EB(KONT)=-.50D0*TK2-.250D0*R4

C(KONT)=0.000

D(KONT)=-.50D0+TK3*RAY((I+1).J.1)-.50D0*FLXR(I)*DS

-.50D0*H(INDEX)*DS*TMPT(J)+(.50D0*TK3+.50D0*H(INDEX)*DS

-.250D0*R3)*RAY(I.J.1)

GO TO 165

2360 CONTINUE

EA(KONT)=-.50D0*TK3-.50D0*H(INDEX)*DS

C(KONT)=-.50D0*TK3

D(KONT)=-.50D0*TK3

D(KONT)=-.50D0*TK3

O(KONT)=-.50D0*TK3

O(KONT)=-.50D0*TK3+.50D0*TK2-.250D0*TK3+.250D0*R3)*RAY(I.J.1)

-.50D0*H(INDEX)*DS*TMPT(J)+(.50D0*TK2-.250D0*R3)*RAY(I.J.1)

O TO 165

VERT SIDE=SEMI-INF* HORIZ CONST FLUX
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1293)
1294) C
                                                                                                                                                                                                       VERT SIDE=SEMI-INF. HORIZ CONST FLUX
```

これからなる。 ありがはない とい

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1295)
12971)
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               c<sup>231</sup>
     1352)
1353)
1354)
1355)
     CHECK WI

CONTINUE

A ROW SET

L=X
IF=1
1368)
CALL TRIDIG(L.EA.EB.C.D)
DO 1607 K=1.X
TEMP(I.K) = D(K)
1371)
1607
CONTINUE
CONTINUE
1372)
160 CONTINUE
1373)
C
JFLG=2
KONT=0
1376)
DO 173 I=1.Y
1378)
DO 174 J=1.Y
1380)
1379
1380)
177
     1360)
                            CHECK WHICH PASS CONTINUE
                                                         A ROW SET UP ON 1ST PASS NOW.
                                                      THE WHOLE GRID 1ST PASS COMPUTED NOW.
                           JFLG=2
KONT=0
D0 173 I=1+Y
D0 174 J=1.X
OLDT(I+J)=RAY(I+J+1)
RAY(I+J+1)=TEMP(I+J)
CONTINUE
CONTINUE
G0 T0 2001
CONTINUE
CONTINUE
CONTINUE
                   174
173
     1381)
1382)
1383)
1384)
1385)
                  A COLUMN SET UP ON 2ND PASS NOW.

1613 CONTINUE
IF=1
L=Y
     1386)
     1388)
```

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1389)
1390)
1391)
1391)
1392)
1608 CONTINUE
1393)
1600 CONTINUE
1394)
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1000 CONTINUE
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1398)
                                  1608 CONTINUE
1600 CONTINUE

WHOLE GRID 2ND PASS COMPLETE NOW.

DO 1609 J=1.X

DO 1610 I=1.Y

RAY(I.J.) = TEMP(I.J)

C CHECK IF FRONT SKIPPED THE NODE

K=RAY(I.J.3)

CC=TPC(K)+(TDEL/2.D0)

DD=TPC(K)-(TDEL/2.D0)

IF((OLDT(I.J).GT.CC).AND.(TEMP(I.J).LT.DD)) GO TO 188

IF((OLDT(I.J).LT.DD).AND.(TEMP(I.J).GT.CC)) GO TO 189

GO TO 1610

188 CONTINUE

C SKIPPED FROM UNFROZEN TO FROZEN

RAY(I.J.1) = CC+(CP(I.J)/CPPC(K))*(RAY(I.J.1)-CC)

GO TO 1610

189 CONTINUE

CC SKIPPED FROM FROZEN TO UNFROZEN

RAY(I.J.1) = DD+(CP(I.J)/CPPC(K))*(RAY(I.J.1)-DD)

1610 CONTINUE

1609 CONTINUE

JFLG=1

KONT=0

Q=Q+1

ITCC=ITCC+1

ITPP=ITPP+1

C
  1398) C
  1408)
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1433)
                                                      WRITE OUTPUT

ITRT IS NO. OF TIME STEPS BEFORE TEMPERATURES PRINTED

ITCC COUNTS FROM 1 TO ITRT

ITPC IS NO. OF TIME STEPS BEFORE ISOTHM IS CALLED

ITPP COUNTS FROM 1 TO ITPC
                                                                TT=Q+DELT
                                                                FIGURE & WRITE ISOTHERM LOCATIONS IF(ITPP-NE-ITPC) GO TO 920 CALL ISOTHM ITPP=0
                                                         WRITE TEMPERATURE DISTRIBUTIONS

IF(ITCC.NE.ITRT) GO TO 888

WRITE(11.893) Q.TT

FORMAT(/,*/,*/1X.**TEMPERATURES AFTER*.14.** TIME STEPS (*.

& F12.3.** HRS):*)

RRX=X/17

JJY=PP*
                                     893
                                                                 JJX=RRX
JJR=X-(17±JJX)
DO 531 JJ=1•JJX
IF(JJX-EQ-0) GO TO 531
  1441)
1442)
1443)
     1444)
1445)
1446)
1447)
1450)
1451)
1452)
1453)
1455)
1457)
                                                                 J2=17*JJ
J1=J2-16
WRITE(11+500) ((RAY(I+J+1)+J=J1+J2)+I=1+Y)
IF(JJ-LT-JJX) WRITE(11+503)
CONTINUE
                                     531
                                   531 CONTINUE
    J2=X
    J1=X-JJR+1
    IF(JJR-EG.0) GO TO 532
    IF(JJX-NE.0) WRITE(11,503)
    OO 532 I=1,7
    WRITE(11,500) (RAY(I,J,1),J=J1,J2)
532 CONTINUE
    ITCC=0
    BBB CONTINUE
                           888
C **
1459)
1460)
14662)
14663)
14663)
14665)
14669)
14772)
14773)
147773)
147773)
147773)
                              C CHECK NO. OF WHOLE TIME STEPS AND GO BACK TO START IF NOT DONE

IF (Q.GE.IMAX) GO TO 172

JFLG=1

KONT=0

GO TO 2002

172 CONTINUE

200 CONTINUE

C WRITE BOUNDARY TYPES INTO RESULT

WRITE (6.404)

404 FORMAT(/,2X,*BOUNDARY CONDITIONS:*)

DO 410 K=1.8

KKK=K

IF (IFLAG(K).EQ.0) GO TO 410

GO TO (400.450.401.451.402.452.403.453).KKK

400 WRITE (6.405)

405 FORMAT(/,1X,*TOP BOUNDARY*)

GO TO 470

401 WRITE (6.406)

406 FORMAT(IX,*LEFT BOUNDARY*)

GO TO 470

407 FORMAT(IX,*BOTTOM BOUNDARY*)
  1480)
1481)
1482)
```

```
GO TO 470

WRITE (6.40B)
FORMATCIX*RIGHT BOUNDARY*)
GO TO 470
FOR MATCIX*RIGHT BOUNDARY*)
GO TO 4600
FORMATCIX*BOTTOP LEFT CORNER*)
GO TO 470
WRITE (6.461)
FORMATCIX*BOTTOM RIGHT CORNER*)
GO TO 470
WRITE (6.462)
FORMATCIX*STOP RIGHT CORNER*)
GO TO 470
CONTINUE
GO TO 470
CONTINUE
GO TO 470
WRITE (6.463)
FORMATCIX*TOP RIGHT CORNER*)
GO TO 470
CONTINUE
GO TO 470
WRITE (6.493)
FORMATCIX*TOP RIGHT CORNER*)
GO TO 410
WRITE (6.493)
FORMATCIX**TOP RIGHT CORNER*)
GO TO 410
WRITE (6.493)
FORMATCIX**
FORMATCIX**
GO TO 410
WRITE (6.493)
FORMATCIX**
GO TO 410
WRITE (6.493)
FORMATCIX**
FORMATCIX**
GO TO 410
WRITE (6.485)
FORMATCIX**
GO TO 410
WRITE (6.486)
FORMATCIX**
GO TO 410
WRITE (6.481)
FORMATCIX**
WRITE THE FINAL TEMPS & THERMAL VALUES
GO TO 510 KK=1.55
GO TO 510 KK=1.55
GO TO 510 KK=1.55
GO TO 510 KK=1.55
GO TO 511 ST.25
GO TO 510 KK=1.55
GO TO 511 ST.25
GO TO 511 ST.25
GO TO 510 KK=1.55
GO
1483)
1484)
1485)
1486)
1487)
                                             408
                                            450
460
 14889)
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463
                                            470
 15001)
15002)
15003)
15005)
15006)
15008)
155078)
155123)
15513
                                            471
492
                                                                                                                                                                                                                                                                                               CONST FLUX ON BOTH SIDES .)
                                           472
493
                                                                                                                                                                                                                                                                           CONVECTIVE ON BOTH SIDES*)
                                           473
494
                                                                                                                                                                                                                                                                                       SEMI-INF ON BOTH SIDES *)
                                           474
484
                                                                                                                                                                                                                                                                                              VERT CONST FLUX--HORIZ CONVECT*)
                                            475
485
                                                                                                                                                                                                                                                                           VERT CONST FLUX -- HORIZ SEMI-INF *)
1516)
                                            476
 1518)
1519)
1520)
                                            486
                                                                                                                                                                                                                                                                                                     VERT CONVECT--HORIZ SEMI-INF*)
                                             477
 1521)
1522)
1523)
                                                                                                                                                                                                                                                                                              VERT CONVECT -- HORIZ CONST FLUX*)
                                            487
                                             478
15245)
152531
15227
155227
155331
155332
1155333
155333
155336
155443
155443
115547
115547
                                            488
                                                                                                                                                                                                                                                                                       VERT SEMI-INF -- HORIZ CONVECT*)
                                            489
                                                                                                                                                                                                                                                                                  VERT SEMI-INF -- HORIZ CONST FLUX*)
                                             491
                                                                                                                                                                                                                                                                                                           CONSTANT TEMPERATURE .)
                                           481
                                            483
                                           416
417
418
410
                                           183
 1548)
1549)
1550)
1551)
1552)
1553)
1555)
                                                                        WRITE THE FINAL TEMPS & THERMAL VALUES DO 510 KK=1.5 CO TO (511.512.513.514.515).KK WRITE(6.893) Q.DD GO TO 516 WRITE(6.911) GO TO 516 WRITE(6.912) GO TO 516 WRITE(6.914) GO TO 516 WRITE(6.914) GO TO 516 WRITE(6.914) GO TO 516 WRITE(6.913) GO TO 518 J2=17.4JJ
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                                           511
                                           513
                                            514
                                            515
                                                                         ÎF(JJX.EQ.0) GO TO 518

J2=17*JJ

J1=J2-16

GO TO (519,520,521,522,523).KK

WRITE(6,500) ((RAY(I,J,1),J=J1,J2),I=1,Y)

GO TO 517

WRITE(6,500) ((TK(I,J),J=J1,J2),I=1,Y)

GO TO 517

WRITE(6,500) ((CP(I,J),J=J1,J2),I=1,Y)

GO TO 517
                                            520
                                             521
```

PERSONAL PROPERTY.

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1577)
1578)
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1586)
1586)
1587)
                                                                                                              WRITE(6.500) ((RO(I.J).J=J1.J2).I=1.Y)
GO TO 517
WRITE(6.502) ((ISTAT(I.J).J=J1.J2).I=1.Y)
IF(JJ.LT.JJX) WRITE(6.503)
CONTINUE
                                                                    522
                                                            523
517
518
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                                                                  SUBROUTINE TRIDIG(N.A.B.C.D)

IMPLICIT INTEGER + 2(I.N)

IMPLICIT DOUBLE PRECISION(A.B.C.D)

DIMENSION A(1).B(1).C(1).D(1)

ADJUST COEFFS FOR B&D FROM ELIMINATING A

DO 10 I=2.N

AB=A(I)/B(I-1)

B(I)=B(I)-AB+C(I-1)

D(I)=D(I)-AB+D(I-1)

CONTINUE

N1=N-1
1643)

1643)

16445)

16447)

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16450)

16550)

16550)

16553)

16553)

16553)

16653)

16662)

16663)

16663)

166667)

16668)
                                                                                                               BACK S
N1=N-1
D(N)=D(N)/B(N)
DO 20 I=1,N1
M=N-I
                                                                                                                 D(M)=(D(M)-C(M)+D(M+1))/B(M)
CONTINUE
RETURN
END
                                                                          **********************************
                                                                                      SUBROUTINE ISOTHM
THIS PROGRAM FINDS ISOTHERMS IN RAY(I.J.1).
IMPLICIT INTEGER 2(A.B.I-L.N.G.V.X.Y.Z)
IMPLICIT DOUBLE PRECISION(C-H.M.O.P.R-V)
DIMENSION EXRY(9.200). EYRY(9.200).COUNT(9)
COMMON/M12I/ X.Y.NISO.Q
```

```
COMMON/M12R/ RAY(80.80.3).TISO(9).DS.DELT
NISO=NUMBER OF ISOTHERMS
COUNT(B)=COUNTER FOR NO. ELTS IN EACH ISOTHERM
TISO(B)=TEMP OF EACH ISOTHERM
LET TISO(1) BE THE HOTTEST ISOTHERM
EXRY(B,K)= ARRAY FOR X-COORDINATES
B INDICATES WHICH ISOTHERM
K=COUNT(B) & INDICATES POSITION OF PT IN ISOTHERM LIST
   1670)
1671) C
1672) C
1673) C
1674) C
1675) C
1676) C
1677) C
1679) C
                                              CHANGE THE NEXT STATEMENT TO AGREE WITH DIMENSION

DO 6 K=1.200

DO 7 B=1.NISO

EXRY(B.K)=0

EYRY(B.K)=0

CONTINUE

CONTINUE

DO 8 K=1.9

CONTINUE

FOR LEFT SEMI-INNF.LET XL=3

FOR BOTTOM SEMI-INNF.LET XL=3

FOR RIGHT SEMI-INNF.LET XR=X-3

FOR RIGHT SEMI-INF.LET XR=X-3

FOR TOP SEMI-INF.LET YT=3. AND COMMENT OUT LOOP 860

IF NOT USING SEMI-INF.XL=1.XR=X-1.YT=2.YB=Y
  16885)
16885)
16885)
168867)
168867)
168997
16997
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16997
16997
16997
16997
                                   6
                                   8
                                                          IF NOT USI

XR=X-1

YB=Y-3

YT=2

DO 100 I=YT+YB

DO 200 J=XL+XR
1701)
1702)
1703)
```

A CANAL MOUNTAIN THE SECOND ASSESSMENT OF

```
GO TO 851
CONTINUE
FYCD=(I-1.)+(RI1-TISO(B))/(RI1-RAY(I,X,1))
EYCO=DS+(EYCD-1.)
COUNT(B)=COUNT(B)+1
K=COUNT(B)
EYRY(B*K)=EYCO
EXRY(B*K)=DS+(X-1.)
CONTPT=CNTPT+1
GO TO 851
CONTINUE
COUNT(B)=COUNT(B)+1
K=COUNT(B)
EYRY(B*K)=DS+(X-1.)
EXRY(B*K)=DS+(X-1.)
CNTPT=CNTPT+1
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
860 LOOP IS FOR TOP ROW
                                                                                                                                                                       1761)
1762)
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1793)
                                                                                                                                                                                                                                                                                                                   XX=X-1
                                                                                                                                                                                                                                                                                                           XX=X-1

I=1

DO 860 J=XL,XK

RJ=RAY(1,J,1)

DO 861 B=1,NIS0

IF((ISO(B).GI.RJ).AND.(IISO(B).LI.RJ1)) GO TO 862

IF((ISO(B).GI.RJ).AND.(IISO(B).GI.RJ1)) GO TO 862

IF(RJ.EQ.TISO(B)) GO TO 863

GO TO 861

CONTINUE

EXCD=(RAY(I.J.1)-TISO(B))/(RAY(I.J.1)-RAY(I.(J.1).1))+J

EXCD=DS*(EXCD-1.)

COUNT(B)=COUNT(B)+1

K=COUNT(B)

EXRY(B.K)=D

CNTPT=CNTPT+1

GO TO 861

CONTINUE

COUNT(B)=COUNT(B)+1

K=COUNT(B)

EXRY(B.K)=D

CNTPT=CNTPT+1

CONTINUE

COUNT(B)=COUNT(B)+1

K=COUNT(B)

EXRY(B.K)=D

CNTPT=CNTPT+1

CONTINUE

CONTINUE

CONTINUE
                                                                                                                                                                                                                                                         862
                                                                                                                                                                           1794)
1795)
                                                                                                                                                                           1796)
1797)
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                                                                                                                                                                        1799) 863
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                                                                                                                                                                                                                                                                                                     TIM=Q*DELT
WRITE(13.96) CNTPT.TIM
FORMAT(1X.0*THE FOLLOWING*,15.0* POINTS REPRESENT TIME=*,

$ F8.4.0* HOURS:*)
FORMAT(1X.5.5)
DO 666 B=1.NISO
L=COUNT(B)
IF(L.EQ.0) GO TO 99
WRITE(13.91) L.B.TISO(B)
FORMAT(1X.216.F6.2)
CONTINUE
DO 95 K=1.L
IF(L.EQ.0.) GO TO 95
WRITE(13.90) (EXRY(B.K).EYRY(B.K))
CONTINUE
FORMAT(1X.2F15.6)
CONTINUE
FORMAT(1X.0*TOTAL NO. POINTS FOUND=*.15)
RETURN
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CPILLUS SECURIC GEAT COUNTRY (WAME)

ROLLING SERVICE CONTROL OF SERVIC
1895)
1896)
1897)
  1898)
1934)
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1947) C INITIALIZE VARIABLES AND ARRAYS
1948) TSRF=-4.670D0
1949) TBTM=4.670D0
1950) TLFT=4.670D0
1951) TRIT=4.670D0
1952) DS=.0050D0
1953) DELT=.00250D0
1954) DI=50.0D0
1955) A=1
1956) IMAX=1200
1957) X=3
1958) Y=25
1959) TDEL=1.0D0
NISO=1
                                                                            X=3
Y=25
TDEL=1.0D0
NISO=1
ITRT=40
ITPC=10
1958;
1959)
1960)
NISO=1
1961)
11RT=40
1962)
1963)
C
1964)
C
1964)
C
1966)
HTIME=IMAX*DELT
1966)
WRITE(1.3)
ITRT.ETIME.IMAX.MTIME
1967)
1968)
S
FORMAT(1X..TEMPERATURES WILL BE PRINTED EVERY..I3.. ITNS../.
1969)
C
1X.. (EVERY..F12.4.. HRS.).././.IX.
1970)
C
MAX NO. ITNS=.F14.. (= .F12.4.. HRS.).

MATERIAL PROPERTIES FOR EACH MATERIAL.

""MATERIAL PROPERTIES FOR EACH MATERIAL.
     1967)
1968)
1970)
1971) C
1971) C
1973) C
1974)
1975) C
1976)
1977)
1978)
1979) C
                                                 SET UP MATERIAL PROPERTIES FOR EACH MATERIAL.

(YOU SET THE VALUES OF K & CP & RO IN MAIN PROGRAM.)

H(A)=0.00

PROPERTIES FOR PHASE CHANGE:

ROPC(1)=917.000

HL(1)=93.000

TPC(1)=0.000

THE FOLLOWING SPECIFIC HEAT SHOULD INCLUDE LATENT HEAT.

YOU CALCULATE IT BY: CP=(CP(FROZ)+CP(UNFROZ))/2 + HL(K)/TDEL

DO 2 K=1.4

CPPC(K)=093.870D0

CONTINUE
     CCC CCC
198834511988567119988567119988507119999111999997811199997811199999
                                               2
                                                SPECIFY THE ISOTHERMS TO BE LOCATED IN SUBROUTINE ISOTHM IN ORDER, WITH TISO(1) THE HOTTEST. TISO(1)=0.0D0
                                                INDICATE THE FLUXES FROM THE SIDES OF THE GRID.

THIS IS USED ONLY FOR THE CONSTANT FLUX BOUND.CONDITION.

DO 11 J=1.X

FLXT(J)=0.D0

FLXB(J)=0.D0

11 CONTINUE

DO 12 I=1.Y

FLXL(I)=0.D0

FLXR(I)=0.D0

12 CONTINUE
                                   C -
1998) 12 CONTINUE
1999) C INDICATE THE TEMPERATURES OUTSIDE THE
2001) C USED FOR SEMI-INF & CONVECT BOUN
2002) TMPT(J)=TSRF
2004) TMPB(J)=TBTM
2005) 13 CONTINUE
2006) DO 14 I=1.Y
2007) TMPL(I)=TLFT
2008) TMPR(I)=TRIT
2009) 14 CONTINUE
2010) C SET UP RAY(I.J.3) INDEX OF MATERIALS
2011) C SET UP RAY(I.J.3) INDEX OF MATERIALS
2012) C THERE ARE A MATERIALS. LET THE SURRO
2013) C MATERIAL BE THE *ATH* MATERIAL TO
2014) C THIS LOOP ASSIGNS THE *ATH* MATERIAL TO
2015) C THIS LOOP ASSIGNS THE *ATH* MATERIAL TO
2016) C THE REMAINING NODES.
2017) DO 10 J=1.X
2018) DO 20 I=1.Y
2019) RAY(I.J.3)=A
2020) 20 CONTINUE
2021) 10 CONTINUE
2022) C SET UP RAY(I.J.2) NODAL LOCATION TYPE
2023) C SET UP RAY(I.J.2)=7
2024) C OONTINUE
2025) C OONTINUE
2027) 40 CONTINUE
2028) C BOTTOM BOUNDARY
2029) CONTINUE
2031) 50 CONTINUE
2033) C LEFT BOUNDARY
2033) C LEFT BOUNDARY
2033) C SET UP CORNER CONDITIONS

58
                                               INDICATE THE TEMPERATURES OUTSIDE THE GRID.

USED FOR SEMI-INF & CONVECT BOUNDARYS

DO 13 J=1.x
TMPT(J)=TSRF
TMPB(J)=TBTM

CONTINUE
DO 14 I=1.y
TMPL(I)=TLFT
TMPR(I)=TRIT

CONTINUE

CONTINUE
                                              SET UP RAY(I.J.3) INDEX OF MATERIALS
THERE ARE A MATERIALS. LET THE SURROUNDING
MATERIAL BE THE *ATH* MATERIAL. START
WITH *1..
THIS LOOP ASSIGNS THE *ATH* MATERIAL TO
THE REMAINING NODES.
DO 10 J=1.x
DO 20 I=1.y
RAY(I.J.3)=A
20 CONTINUE
```

ASSESSED AND ASSESSED OF THE PARTY AND ASSESSED

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2041) C --DON*T CHANGE THE "RAY(I,J,2)" STATEMENTS.
2042) RAY(I,1,2)=1
2043) RAY(Y,1,2)=2
2044) RAY(Y,0,2)=3
2045) RAY(1,0,2)=4
2046) C TOP LEFT CORNER
2047) RAY(1,1,1)=TSRF
2049) C RRNR(1)=1
2049) C RRNR(2)=7
2051) RAY(Y,1,1)=TBTM
2052) C RNR(3)=7
RAY(Y,1,1)=TBTM
2053) RAY(Y,0,1)=TBTM
     2047)
2048)
2049)
2050)
2051)
2053)
2055)
C
                                                                                                                                       RAY(Y, X, 1) = TBTM
TOP KIGHT CORNER
     2055)
2056)
2057)
2058)
2060)
2061)
2062)
2063)
                                                                                                                                         KRNR (4)=1
                                                                                                                                       RAY(1.X.1)=TSRF
                                                                                    INTERIOR NODES GIVEN FLUCTUATING TEMP
DO NOT CHANGE THE 80.90 LOOP.
YY=Y-1
                                                                                                                                      XX=X-1

DO 80 J=2.XX

DO 90 I=2.YY

RAY(I.J.2)=5

CONTINUE
        2064)
2065)
    2065) RAY(I,J,2)=5
2066) 90 CONTINUE
2067) 80 CONTINUE
2069) C CONSTANT INTERIOR NODES.
2070) C INDICATE RAY(I,J,2)=6 IF RELEVANT
2071) C
2072) C IF RELEVENT. SPECIFY SPECIAL RAY(I, 2073) C
2073) C TO BE USED WITH SEMI-INF BOUND(2075) XX=X-1
2075) YY=Y-1
2076) DO 210 J=2.XX
RAY(YY,J,2)=12
2077) RAY(YY,J,2)=12
                                                                                   IF RELEVENT. SPECIFY SPECIAL RAY(I.J.2) CONDITIONS

TO BE USED WITH SEMI-INF BOUNDARY.

XX=X-1
YY=Y-1
DO 210 J=2.XX
RAY(YY.J.2)=12
210 CONTINUE
RAY(YY.2.2)=15
RAY(YY.1.2)=23
RAY(Y.2.2)=19
RAY(YY.XX.2)=15
RAY(YY.XX.2)=15
RAY(YY.XX.2)=21
       2078) 210
2079) C
2080)
       2081) C
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                                                                                      SET UP RAY(I.J.1) PRESENT NODAL TEMP
                                                                      2094) RAY(II.10.1)=TLFT
2096) 120 CONTINUE
2097) XXXX-1
2098) DO 140 I=2.YY
2100) DO 140 I=2.YY
2101) C COMMENT OUT THE INITIAL TEMP DIST THAT YOU DON'T WANT:
2102) C NODES ASSIGNED LINEAR TEMP DIST VERTICALLY
2103) C NODES ASSIGNED LINEAR TEMP DIST VERTICALLY
2104) C NODES ASSIGNED UNIFORM TEMP (INDICATE THE TEMP:)
2105) RAY(I.J.1)=4.670D0
2106) 140 CONTINUE
2107) 130 CONTINUE
2108) C INTERIOR NODES
2110) C INTERIOR NODES
2110) C INTERIOR NODES
2111) C IF THEY ARE DIFFERENT THAN ASSIGNED IN THE 130.140 LOOP.
2112) C
2113) C
2113) C
2114) C DO NOT CHANGE THE FOLLOWING 17 STATEMENTS.
2115) WRITE(5.131) DS.DELT.DI.TDEL
2116) 131 FORMAT(IX.49F10.5)
2117) WRITE(5.132) A.X.Y.NISO.ITRT.HAX.ITPC
2118) 132 FORMAT(IX.715)
2119) WRITE(5.134)(TISO(B).B=1.NISO)
2121) 134 FORMAT(IX.755)
2122) WRITE(5.133) (H(L).L=1.A)
2122) WRITE(5.133) (H(L).L=1.A)
2123) WRITE(5.131) (RAY(I.J.3).J=1.X).I=1.Y)
2124) 155 FORMAT(IX.9F10.5)
2125) WRITE(5.151) ((RAY(I.J.3).J=1.X).I=1.Y)
2126) WRITE(5.151) ((RAY(I.J.3).J=1.X).I=1.Y)
2127) WRITE(5.151) (FLXE(J).J=1.X)
2131) WRITE(5.151) (FLXE(J).J=1.X)
2131) WRITE(5.151) (FLXE(J).J=1.X)
2131) WRITE(5.151) (FLXE(J).J=1.X)
2132) WRITE(5.151) (FLXE(J).J=1.X)
2134) WRITE(5.151) (FLXE(J).J=1.X)
2134) WRITE(5.151) (FLXE(J).J=1.X)
```

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2135) WRITE(5.151) (TMPT(J).J=1.X)
2136) WRITE(5.151) (TMPB(J).J=1.X)
2137) WRITE(5.151) (TMPL(I).I=1.Y)
2138) WRITE(5.151) (TMPR(I).I=1.Y)
2139) RETURN
2140) END
2141) C
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Action to the second in the second account to the second s

Data file ADPDAT printed by subroutine ADDATA for the Neumann problem

The second of the second secon

	0000	22.25 5.00 5.00 6.00 6.00 6.00 6.00 6.00 6.0	444 646 777	00.0	• •	4.67	1.67
	0000	25.000	**** **** *****	0.00	00.0	4.67	4.67
	0000	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	444 •••• •••• •••• ••••	00.0	0.0	4.67	4.67
	0000	27.00 29.00 59.00	44.67 76.64 76.64 76.64	ē	0 0 0	4.67	4.67
	0000	27.00 29.00 27.00	**** **** **** ****	0	00.0	4.67	4.67
	0000	29.00 27.00 29.00	4444 2000 7777	00.0	00.	4.67	4.67
	0000	27.00	444 6667	0.00	00.00	4.67	4.67
	0000	27.00 29.00 5.00	4444 6664 7777	0.00	00.0	4.67	4.67
	0000	29.00 27.00 29.00	4444 767 767	0.00	0.00	4.67	4.67
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•			61	-			

File ADPOUT (Lists the initial data in readable form and the final values from the run [printed by ADIPC].)

```
DATA FOR THIS RUN OF ADIPC:
                             DELT DI 0.00250 50.00000
         05
                                                                             TDEL
1.00000
                                Y NISO ITRT IMAX
25 1 40 1200
    ROPC(K) CPPC(K) HL(K)
917-00000 93-87000 93-00000
 TISO(B).B=1.NISO:
      H(K)
0.00000
 KRNR(1) = 1 KRNR(2) = 7 KRNR(3) = 7 KRNR(4) = 1
      FLXT(J)
0.00
0.00
0.00
                            FLXB(J)
0.00
0.00
0.00
                                                  TMPT(J)
-4.67
-4.67
-4.67
                                                                          TMPB(J)
4.67
4.67
4.67
                            FLXR(I)
0.000
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     FLXL(I)
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                                                   TMPL(I)
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                                                                           TMPR(I)
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RAY(I+J+1) TI
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                               TEMPERATURES AT THE START:
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RAY(1.J.2) NO
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27.00 5.00
                             NODAL LOCATION TYPE

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NODAL MATERIAL TYPE

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                                                                     1.00
  TEMPERATURES WILL BE PRINTED EVERY 40 TIME STEPS. ISOTHERMS WILL BE LOCATED EVERY 10 TIME STEPS.
      BOUNDARY CONDITIONS:
 TOP BOUNDARY

TOP LEFT CORNER

CONSTANT TEMPERATURE
CONSTANT TEMPERATURE
LEFT BOUNDARY

CONSTANT HEAT FLUX
BOTTOM LEFT CORNER VERT CONST FLUX--HORIZ SEMI-INF
BOTTOM BOUNDARY SEMI-INFINITE
BOTTOM RIGHT CORNERVERT CONST FLUX--HORIZ SEMI-INF
RIGHT BOUNDARY CONSTANT HEAT FLUX
TOP RIGHT CORNER

CONSTANT TEMPERATURE
  NUMBER OF TIME STEPS= 1200
3.000 HRS):
```

CANADA OF CANADA

Caron exercised

FINAL 21 20	FINAL RO (1.J): 917.00 917.00 917.00 917.00 917.00 917.00 917.00 917.00 917.00 917.00 917.00 9917.00 917.00 9917.00 9918.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20 998.20
FINAL CP(I,J): 0.58 0.58 0.58 0.58 0.58 0.58 0.58 0.5	FINAL ISTAT(I.J): 11 11 11 11 11 11 11 11 11 11 11 11 11

BANNERS REPRESENTATION CALLES CONTINUES FOR THE PROPERTY SECTIONS FOR THE PARTY OF THE PARTY OF

(Contains the values from the end of the run [calculated by ADIPC] and could be used to start the program again to run for more time.) Data file ADPNDT

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	0000	22.25 29.00 59.00	0 0 0 0 0 0 0 0 0 0 0 0 0	00.0	00.0	4.67	4.67
	0000 0000	27.00 27.00 27.00 00 00	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 • 0	0.0	4.67	4.67
	4444 0000 0000	24.00	0.18.4 0.84 0.44	0.00	0.00	4.67	4.67
	4444 0000 0000	27.00 29.00 5.00	0.124 0.034 1.450	00.00	0.00	4.67	4.67
	4111 0000 0000	27.00 29.00 27.00	0184 0450 7450	00.00	0.00	4.67	4.67
	4444 0000 0000	29 27 29 00 29 00 00	11.84 00.46 00.46	00.0	0.00	4.67	4.67
	4444 0000 0000	5.00 27.00 29.00 5.00	11 K &	00.0	0.00	4.67	4.67
	0000 0000	27.00	11.04 0000 0000	0.00	0.0	4.67	4.67
	4444 0000 0000	29.00	3415 340 340 317 317 317 317 317 317 317 317 317 317	00*0	0 • 0	4.67	4.67
	0000 0000	27.00 29.00 5.00	-2.82 1.07 3.79	0	000	9.4	4.67
	0000		0000W4 000000 0000000	0,0	000	94	4.67
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50.0000 40.120 93.0000	0000		1002W4	Ö	200	94	4.67
00250 87000	0000		1402W4 2400W4 240044	0000	14000	90	40
	9000		0400W4 0400W4 040W4401 040044	0000	2004	9.04	99
000		20000			9009	994	99

Sample of isotherm locations output from ADIPC for the Neuman problem (file POINT1)

```
3 POINTS REPRESENT TIME= 0.0250 HOURS:
0.005000 0.004632
0.010000 0.004632
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.0500 HOURS:
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.0750 HOURS: 0.000000 0.004830
THE FOLLOWING 3 POINTS REPRESENT TIME=
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.1250 HOURS: 0.000000 0.005197
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.1500 HOURS: 0.000000 0.005661
0.005000 0.005661
0.010000 0.005661
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.1750 HOURS:
3 0.000000 0.006161
3 POINTS REPRESENT TIME = 0.2000 HOURS:
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.2250 HOURS: 0.000000 0.007223
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.2500 HOURS:
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.2750 HOURS: 0.00000 0.009413 0.005000 0.009413 0.005000 0.009413
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.3000 HOURS:
                               REPRESENT TIME= 0.3250 HOURS:
                          0.009618
0.009618
0.009618
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.3500 HOURS:
0.005000 0.009724
0.010000 0.009724
THE FOLLOWING 3 POINTS REPRESENT TIME= 0.3750 HOURS:
THE FOLLOWING 3 POINTS REPRESENT TIME = 0.4250 HOURS: 3 0.00000 0.010158 0.010000 0.010158
```

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Albert, Mary Remley

Computer models for two-dimensional transient heat conduction / by Mary Remley Albert. Hanover, N.H.: Cold Regions Research and Engineering Laboratory; Springfield, Va.: available from National Technical Information Service, 1983.

iv, 74 p., illus.; 28 cm. (CRREL Report 83-12.) Prepared for Office of the Chief of Engineers by Corps of Engineers, U.S. Army Cold Regions Research and Engineering Laboratory under DA Project 4A762730 AT42.

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12-83.